

1 **Comparison between the non orographic gravity wave**
2 **drag schemes used in global climate models to simulate**
3 **a quasi-biennial oscillation and constant level balloons**

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14 **Key Points:**

- 15 • **Temptative, but somethink like:**
- 16 • The non-orographic parameterization tuned to produce a realistic tropical quasi-
17 biennial oscillation in global climate models are used to predict in-situ observa-
18 tions.
- 19 • Parameterized gravity waves needed in large-scale models have realistic amplitudes
20 in the tropical lower stratosphere.
- 21 • Day-to-day variations of the estimated gravity wave momentum fluxes correlate
22 in some cases with observations, except when launching level are near the tropopause.
- 23 • The probability density distribution of the parameterized momentum fluxes are
24 better reproduced when the schemes are not related to their convective sources
25 and/or the launching level in the lower and middle troposphere.

26 **Version 0, date: October 31, 2023**

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Abstract

Gravity Waves (GWs) parameterization schemes in 14 General Circulation Models (GCMs) participating to the Quasi-Biennial Oscillation initiative (QBOi) are directly compared to Strateole-2 balloon observations made in the lower tropical stratosphere from November 2019 to February 2020 (phase 1) and from October 2021 and January 2022 (phase 2). The parameterizations used span the 3 state of the arts techniques used in GCMs to represent subgrid scale non-orographic GWs, the two globally spectral techniques developed by (Hines, 1997) and (Warner & McIntyre, 1999) respectively and the "multiwaves" approaches following (Lindzen, 1981). The input meteorological fields necessary to run the parameterizations offline are extracted from the ERA5 reanalysis and correspond to the instantaneous meteorological conditions found underneath the balloons. In general, the amplitudes are in fair agreement between measurements of the momentum fluxes due to waves with periods less than 1 hr and the parameterizations. The correlation of the daily values between the observations and the results of the parameterization can be around 0.4, which is statistically elevated considering that we analyse around 1200 days of data and sometime good considering that the parameterizations have not been tuned: the schemes used are just the standard ones that help producing a Quasi-Biennial Oscillation in the corresponding model. These correlations nevertheless vary considerably between schemes and depend little on their formulation (globally spectral versus multiwaves for instance). We therefore attribute it to dynamical filtering all schemes taking good care of it, whereas only few relate the gravity waves to their sources. Except for one model, significant correlations are mostly found for eastward propagating waves, which may be due to the fact that during both Strateole 2 phases the QBO phase is easterly at the altitude of the balloon flights. On the other hand, statistical properties, like pdf of momentum fluxes seem better represented in spectral schemes with constant sources than in schemes ("spectral" or "multiwaves") that relate GWs to their convective sources.

Plain Language Summary

In most large-scale atmospheric models, gravity wave parameterizations are based on well understood but simplified theories and parameters which are keyed to reduce systematic errors on the planetary scale winds. In the equatorial regions, the most challenging errors concern the Quasi Biennial Oscillation. Although it has never been verified directly, it is expected that the parameterizations tuned this way should transport a realistic amount of momentum flux in both the eastward and westward directions when compared to direct observations. Here we show that it is the case, to a certain extent, using constant-level balloon observations at 20 km altitude. The method consists in comparing directly, each day and at the location of the balloon the measured momentum fluxes and the estimations from the gravity wave parameterizations used in the global models that participate to the Quasi-Biennial Oscillation initiative and when using observed values of the large-scale meteorological conditions of wind, temperature, precipitation, and diabatic heating.

1 Introduction

It is well known that the large scale circulation in the middle atmosphere is in good part driven by gravity waves (GWs) that propagate in the stratosphere (Andrews et al., 1987). These waves carry horizontal momentum vertically and interact with the large scale flow when they break. The horizontal scale of these waves can be quite short, much shorter than the horizontal scale of General Circulation Models (GCMs) so they need to be parameterized (Alexander & Dunkerton, 1999). In the tropics, the convective GWs are believed to dominate largely (Fovell et al., 1992; Alexander et al., 2000; Lane & Moncrieff, 2008), they contribute significantly to the forcing of the Quasi-Biennial Oscillation (QBO), a near 28-month oscillation of the zonal mean zonal winds that occurs in the lower part of the equatorial stratosphere (Baldwin et al., 2001). For these reasons, the parameterization of convective GWs is necessary for most GCMs to explicitly realize the QBO.

Although gravity wave parameterizations are now used in many models with success including in the tropics (Anstey et al., 2016; Scinocca, 2003; Orr et al., 2010; Christiansen et al., 2016) (Serva et al., 2018)(Beres et al., 2005; Song & Chun, 2005; Lott & Guez, 2013; Bushell et al., 2015), their validation using direct in situ observations remains a challenge. There exist observations of GWs using global satellite observations (Geller et al., 2013) but the GWs identified this way still have quite large horizontal scales, and some important quantities like the Momentum Fluxes (MFs) are often deduced indirectly, for instance from temperature measurements using polarization relations (Alexander et al., 2010; Ern et al., 2014). For these two reasons, in situ observations are essential, and the most precise ones are provided by constant-level long-duration balloons, like those made in the Antarctic region during Strateole-Vorcore (Hertzog, 2007) and Concordiasi (Rabier et al., 2010), or in the deep tropics during PreConcordiasi (Jewtoukoff et al., 2013) and more recently Strateole 2 (Haase et al., 2018). Among many important results, these balloon observations have shown that the momentum flux entering in the stratosphere is extremely intermittent (Hertzog et al., 2012). This intermittency implies that the mean momentum flux is mostly transported by few large-amplitude waves that potentially break at lower altitudes than when the GW field is more uniform. This property, when reproduced by a parameterization (de la Cámara et al., 2014; Kang et al., 2017; Alexander et al., 2021), can help reduce systematic errors in the midlatitudes, for instance on the timing of the final warming in the Southern Hemisphere polar stratosphere (de la Cámara et al., 2016), or on the QBO (Lott, Guez, & Maury, 2012). Balloon observations have also been used to characterize the dynamical filtering by the large scale winds (Plougonven et al., 2017), and to validate the average statistical properties of the GW momentum flux predicted offline using reanalysis data (Kang et al., 2017; Alexander et al., 2021). However, the evaluation of parameterizations using balloon observations have often been quite indirect so far, with the common belief that the best a parameterization can do is to reproduce a realistic statistical behaviour (Jewtoukoff et al., 2015; Kang et al., 2017; Alexander et al., 2021). and there are few good reasons for that. For example, parameterizations are often based on simplified quasi-linear wave theory, they assume spectral distributions that are loosely constrained, and they ignore lateral propagation almost entirely (some attempt to include it can be found in Amemiya and Sato (2016)). Nevertheless, some factors could mitigate these weaknesses. One is that in most parameterizations the wave amplitude is systematically limited by a breaking criterion that encapsulates nonlinear effects. An other is that some parameterizations explicitly relate launched waves to sources, and there is constant effort to improve the realism of the convective ones (Liu et al., 2022). Also, observations systematically suggest that dynamical filtering by the large scale wind is extremely strong for upward propagating GWs (Plougonven et al., 2017), and this central property is represented in most GW parameterizations. For all these reasons, it may well be that GW parameterizations keyed to the large scale conditions found at a given place and time gives MFs that can be directly compared to the MFs measured by a balloon at the same place.

122 Based on the relative success of the offline calculations done in the past using re-
 123 analysis data (Jewtoukoff et al., 2015; Kang et al., 2017; Alexander et al., 2021), (Lott
 124 et al., 2023) have shown that such a direct comparison gives result of interest. The first
 125 is that a state of the art convective gravity wave drag scheme, predicts momentum fluxes
 126 in the low equatorial stratosphere which amplitude can be directly compared with those
 127 measured during the Strateole-2 balloon flights. It gives a direct in-situ observational
 128 confirmation that the theories and modelling of the QBOs developed over the last 50 years
 129 were in good part correct about the significance of the GWs to the QBO forcing. Also
 130 interesting, the comparison showed a good level of correlation between the day to day
 131 variability in momentum fluxes between measured and observed fluxes, a correlation that
 132 is much better for waves carrying momentum fluxes in the eastward direction than in
 133 the westward direction. It was suggested that such a good correlation was due to the
 134 fact that the (Lott & Guez, 2013) scheme analysed relate the gravity waves to their con-
 135 vective sources (not all schemes do) and that the GWs experience significant dynami-
 136 cal filtering in the middle troposphere and lower stratosphere. (Lott et al., 2023) nev-
 137 ertheless revealed that a scheme that relates gravity waves to convection only somehow
 138 failed in predicting the right statistical behaviour of the momentum fluxes, the pdfs show-
 139 ing a long tail for low values of the MFs, suggeting missing processes like lateral prop-
 140 agation or the presence of a background of waves which origin remains a challenge to pre-
 141 dict.

142 The purpose of this paper is to continue such a direct comparison using the most
 143 recent observations and near all the gravity wave parameterization schemes used by the
 144 modelling groups participating to the Quasi-Biennial Oscillation initiative (QBOi, (Butchart
 145 et al., 2018)). We will follow for that (Lott et al., 2023) and use the 8 balloons of the
 146 first phase of the Strateole 2 campaign that flew in the lower tropical stratosphere be-
 147 tween November 2019 and February 2020 and extent it to the 15 balloons that flew more
 148 than one day during the second phase of the Strateole 2 campaign, between October 2021
 149 and January 2022. For each of the flights and each time, we have identified the grid point
 150 in the ERA5 reanalysis (Hersbach et al., 2020) that is the nearest and used the verti-
 151 cal profiles of wind and temperature as well as the surface value of precipitation to em-
 152 ulate the parameterization of GWs used in the global models that participated to the
 153 Quasi-Biennial Oscillation initiative (QBOi). We also extract from the analysis and short
 154 range forecast, diabatic heatings and the cloud base and top altitudes needed in some
 155 schemes to predict gravity waves.

156 2 Data and method

157 2.1 Parameterizations of non orographic gravity wave schemes

158 The parameterizations schemes used to predict orographic gravity waves belongs
 159 to two well separated families, dating back from the 1980's when it becomes evident that
 160 a good simulation of the middle atmosphere by global atmospheric models could not
 161 be done without taking into account non-orographic gravity waves. The first family roots
 162 in the formulation by (Lindzen, 1981), where the gravity wave field field is represented
 163 by gravity waves monochromatic in the horizontal space and time. It was extended to
 164 treat a large ensemble of waves by (Alexander & Dunkerton, 1999) making the assump-
 165 tion that the breaking of each waves could be made independent from the others. An
 166 advantage of such schemes is that it roots in linear theories where sources like convec-
 167 tion and/or fronts can be introduced using closed form theories (Beres et al., 2005; Song
 168 & Chun, 2005; Richter et al., 2010; Lott & Guez, 2013; de la Cámara & Lott, 2015). In
 169 the following we will refer to such schemes as "multiwave", they are expensive because
 170 they request a large amount of waves to represent well a realistic wave field, but this limit
 171 can easily be circumvented by using stochastic approaches (Eckermann, 2011; Lott, Plougonven,
 172 & Vanneste, 2012). As an alternative, but also to better represent breaking, many cen-
 173 ters developed globally spectral schemes. These schemes uses the observationnal fact that

	p_l	F_{LT}	$2\pi/m_*$	C_{\min}
CMAM	100hPa	1.3mPa	1km	0.25 m/s
IFS	450hPa	5mPa	3km	0.5 m/s
ECEarth	450hPa	3.75mPa	2km	0.25 m/s
UKMO	1000hPa	Precip	4.3km	?

Table 1. WMI Parameters changing between CMAM, IFS, ECEarth, and UKMO. UKMO is shown distinctly because it is based on (Warner & McIntyre, 1999) simplified version of WMI rather than on (Scinocca, 2003)’s and realte launched MF to precipitations.

174 in observations the vertical shape of the spectra have a quite universal character. In the
 175 early 1990’s (Hines, 1991) developed a theory where GW breaking is represented by im-
 176 posing an upper limit to the range of vertical wavenumber, the limit being calculated
 177 according to the large scale wind and including a Doppler spreading by the other grav-
 178 ity waves (see also (Hines, 1997)). The scheme has been implemented with success in var-
 179 ious GCMs (see for instance (Manzini et al., 1997)), and will be referred to as ”HDS”
 180 in the following. As an alternative, the theory in (Warner & McIntyre, 1996) imposes
 181 gravity wave saturation according to an empirical spectra but treat vertical changes in
 182 the spectra following GWs propagation invariant character. The theory has been sim-
 183 plified and/or optimized to permit implementation, for instance in the UKMO model (Warner
 184 & McIntyre, 1999; Scaife et al., 2002) and in the CMAM model (Scinocca, 2003) respec-
 185 tively, and will be refered to has ”WMI” in the following. To a certain extent, the spec-
 186 tral schemes can also take into account the relation with sources, for instance the HDS
 187 scheme has been related to fronts in (Charron & Manzini, 2002), and the UKMO ver-
 188 sion of the WMI scheme to precipitations in (Bushell et al., 2015).

189 In the present paper, we are going to compare the GWs schemes used in 14 of the
 190 models that participate to QBOi, all belonging to one of the three type of schemes de-
 191 scribed above (WMI, HDS, and Multiwave), all the multiwave schemes relating GWs to
 192 their convective sources, only one of the spectral scheme doing so, the UKMOgws WMI
 193 scheme in (Bushell et al., 2015), its results will be discussed with the source-related mul-
 194 tiwave scheme.

195 Among the 14 models, three use the (Scinocca, 2003)’s version of the WMI: CMAM,
 196 IFS and ECEarth. Their version for QBOi and further detailed in (Anstey et al., 2016;
 197 Orr et al., 2010; Davini et al., 2017), they essentially differ by four parameters, the launch
 198 level pressure p_l , the launched momentum flux F_{LT} , the characteristic vertical wavenum-
 199 ber m_* and a minimum intrinsic phase speed in the launched spectra Note that for EC-
 200 Earth the exact value of the parameters are from J. Garcia Serrano (private communi-
 201 cation).

202 Still among the 14 models, 5 uses the (Hines, 1997)’s parameterization schemes pre-
 203 sented in (Manzini et al., 1997). Between the five models, and in the Hines scheme, only
 204 changes between models the launching level p_l , the root mean square of the horizontal
 205 wind variability due to GWs at launch level σ , as well as an effective horizontal wavenum-
 206 ber K^* . There are also more numerical parameters that eventually changes, a minimum
 207 value for the the cutoff vertical wavenumber m_{\min} , and two parameters that control smooth-
 208 ing in the vertical of the GWs root mean square variance and cut-off vertical wavenum-
 209 ber, the coefficient C_{SMO} and the number of time the smoothing is applied N_{SMO} .

	p_l	σ_s	$2\pi/K^*$	m_{\min}	C_{sno}	N_{sno}
ECham5	450hPa	1.	125km	0	2	5
MIROC	650hPa	0.95	250km	$6.5 \cdot 10^{-5}$	2	2
MPIM	650hPa	1.2	125km	0	2	2
MRI-ESM	700hPa	1.9	1250km	$3.3 \cdot 10^{-4}$	4	2
EMAC	650hPa	1.	125km	0	2	2

Table 2. HDS Parameters changing between ECHam5, MIROC, MPIM, MRI-ESM, and EMAC.

	p_l	Phase Speed	Δz	Source
LMDz	500hPa	-30m/s<Intrinsic<30m/s	1km	Precip
Yonsei	900hPa-100hPa	-100m/s<Absolute<100m/s	100m-15km	Heating rate
WACCM	?			Heating rate

Table 3. Some parameters changing between LMDz, YONSEI and WACCM, for information only the schemes being extremely distinct

210 Finally the last 4 schemes we consider all links GWs to sources (convection or pre-
211 cipitation), 3 are multiwaves and have been developed independently one from the oth-
212 ers, LMDz, YONSEI, and WACCM and 1 uses the ultra simple version of the WMI schemes
213 presented in (Warner & McIntyre, 1999; Bushell et al., 2015) rather than the (Scinocca,
214 2003)’s version. The differences between these schemes are numerous it is impossible to
215 details them, the reader is referred to the corresponding papers. The most salient dif-
216 ferences are in the source term. In LMDz is made the choice to relate the launched MF
217 to square precipitation P_r^2 consistent with linear theory before breaking (Lott & Guez,
218 2013) whereas in (Bushell et al., 2015) it is related to $\sqrt{P_r}$. In the Yonsei’s scheme (Song
219 & Chun, 2005; Choi & Chun, 2011), the launched momentum flux is directly related to
220 convective heating distributed in the vertical between the cloud bottom and cloud top,
221 the launch altitude being at the cloud top. In this case the launching level can vary be-
222 tween $2km$ and $15km$ typically and the depth of the heating between $100m$ and $15km$.
223 Also, the absolute phase speed cover the ranges $100m/s < C_{abs} < 100m/s$.

224 2.2 Offline parameterization runs

225 To activate the schemes in offline mode we will use ERA-5 hourly data of precip-
226 itation and 3-hourly data of winds, surface pressure, temperature, cloud liquid and ice
227 water content at $1^\circ \times 1^\circ$ horizontal grid to mimic a large scale climate model resolution.
228 Winds, surface pressure, temperature, and water contents are then linearly interpolated
229 on 1hr time step to be synchronised with precipitation. In the vertical we use data at
230 67 model levels, taking one every two ERA5 levels, to speed up calculations. To estimate
231 convective heating rates vertical profiles, we follow (Fueglistaler et al., 2009) and eval-
232 uate diabatic heating using ERA5 hourly data from short range forecast and as a resid-
233 ual between the parameterized temperature tendency and the radiative heatings (long-
234 wave plus shortwave). When needed, we also evaluate the cloud bottom height and top
235 height using the cloud water content (liquid+ice) from ERA5 reanalysis.

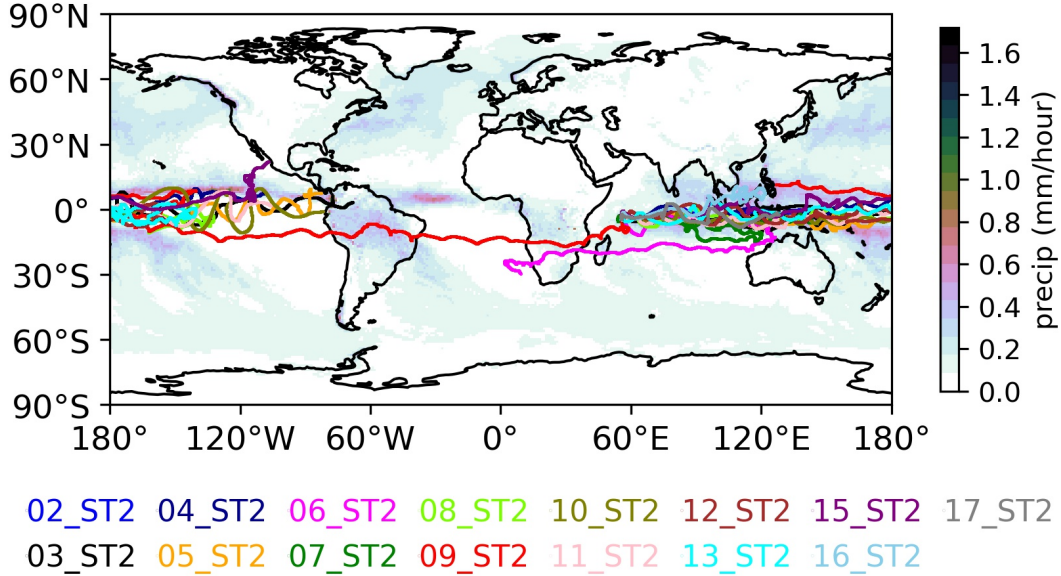


Figure 1. Strateole 2, Phase 2 balloon trajectories taking place between October 2021 and January 2022. Shading presents the precipitation field from ERA5 averaged over the period.

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2.3 Strateole 2 balloon observations

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The in situ observations we use are from the 8 balloons of the first phase of the Strateole 2 campaign that flew in the lower tropical stratosphere between November 2019 and February 2020 and from the 15 balloons that flew more than one day during the second phase of the Strateole 2 campaign, between October 2021 and January 2022. The trajectories during phase 2 are shown in Figure 1, superimposed upon the averaged precipitation (the same Figure but for phase 1 is in (Lott et al., 2023)). In the MFs calculated from observations Corcos et al. (2021) distinguish the waves with short periods (1hr-15mn) from the waves with period up to one day (1d-15mn). They also distinguish the eastward waves giving positive MF in the zonal direction from the westward waves giving negative MF. As shows Fig. ??, it is coincidental that the phase 1 flights took place during the 2nd documented QBO disruption (Anstey et al., 2021), but the fact that the measurements are below the altitude at which the disruption manifests makes us believe that our comparison between gravity wave MFs over the period is not much affected by the disruption (beyond the fact that the disruption potentially affects the large scale winds, which is something that translates well in the parameterization). It is also important to notice that during phase 2, the large scale winds conditions in the lower stratosphere are almost as during phase one at balloons altitudes (end of eastward QBO phase).

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In the following we will compare the momentum fluxes derived from the balloon data, emphasize the intrinsic frequencies that the scheme represents (the intrinsic periods below 1hr) and consider the ERA5 data at the points that are the nearest from the balloon. The prediction is then made every hour and averaged over the day, again because this is the time scale needed for our scheme to sample realistically a GW field, and also because it takes around a day for a balloon flight to cover about a model grid-scale. We will discuss sensitivities to these choices in the first paragraph of the conclusion.

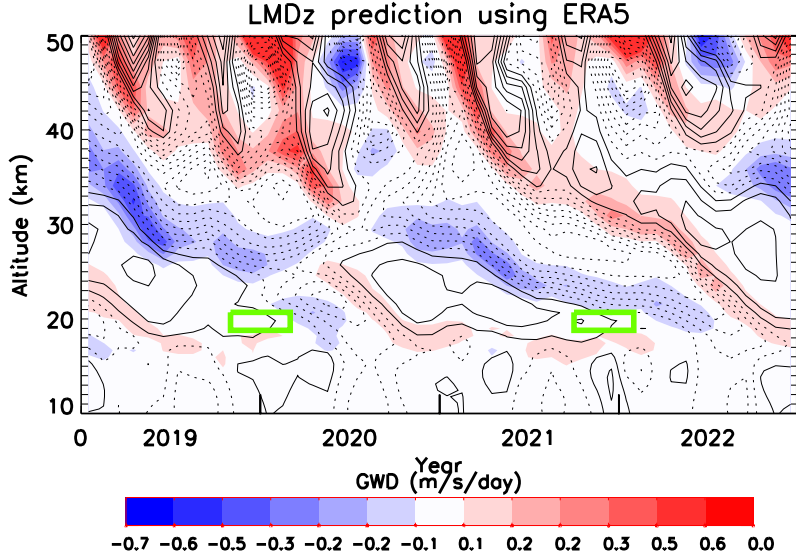


Figure 2. Time vertical sections of the zonal mean zonal wind (CI=10m/s, negative values are dashed and non-orographic gravity wave tendency averaged over the Equatorial band ($-6S - +6N$). Input data are from ERA5 reanalysis and GWs prediction from the LMDz scheme. The 2 green boxes indicate schematically the altitude and time ranges of the Strateole 2 phase 1 and 2 flights considered in this study.

3 Results

THIS SECTION IS UNDER PROGRESS, ONLY THE FIGURES SHOULD BE CONSULTED

Figure 3 shows time series of daily values of momentum fluxes predicted by the parameterizations and measured during balloon flights 2 from strateole 2 phase 1. This is also the flight shown in Fig. 3 in (Lott et al., 2023), and where was also shown the time series of daily precipitation and zonal wind at flight altitude. The top panel is for the WMI based schemes, the middle panel for the HDS schemes and the bottom panels for the schemes relating the GWs fluxes to their sources. In all panels the black curves are for the daily observations. For clarity we present results for the Eastward and westward MF only. Overall ones sees that the schemes predict momentum flux values that somehow compare with the observed one, at least in terms of amplitude. In the globally spectral schemes (upper and middle panels), we also see that the parameterizations predict peaks in MF which duration and amplitudes looks reasonable compared to the observed ones. More specifically, and in the eastward direction there are pronounced peaks after day sixty, peaks that (Lott et al., 2023) related to the fact that after this date the zonal wind at balloon altitude is westward a situation that favors eastward waves. Of course in terms of amplitude there are variabilities between the models schemes, but the qualitative behaviours look reasonable. There is nevertheless one exception, the CMAM scheme predicts long lasting plateau, maybe for this scheme the fact that the launching altitude is quite high, mitigate the dynamical filtering between the launching level and the balloon flight.

The fact that the parameterizations estimates fluxes of about the right amplitude is summarized in Figure 5, where the average of the fluxes over all strateole flights that last than a month (18 flight in total) are shown. It confirms systematically that the offline estimations are quite good on average and in the zonal direction, for the eastward

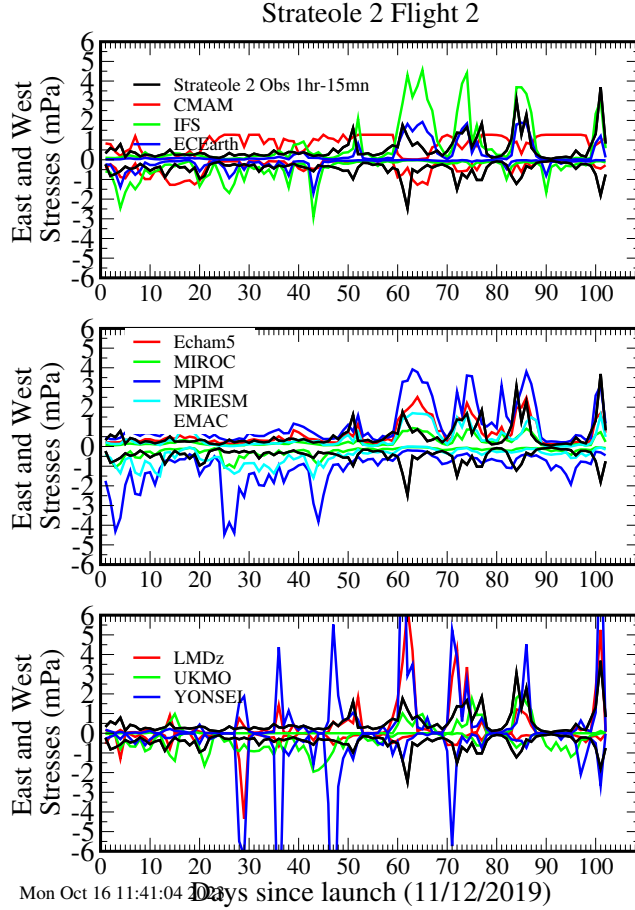


Figure 3. Comparison between daily averaged values of the eastward and westward MFs measured by the balloons during Strateole 2 Flight 1 and estimated by the GW schemes at the balloon location and altitude. Colored curves are for the GW schemes predictions using ERA5 and from different models, black curves are for the observed MFs due to the 15mn-1hr GWs.

288 and westward components again. The Figure 6 group the models averaging the eastward
 289 and westward fluxes over all the balloon flights, confirming again that the parameter-
 290 izations used fall around the observed values. There is varibilities between the models,
 291 but there is no tendency for the models to vcollectively need much more or much less
 292 gravity wave stress than needed.

293 The curves in Figure 3 and Figure 4 also suggest that observations and offline es-
 294 timations evolve quite similarly day after day, both measured and parameterized MFs
 295 being sensitive to precipitation and dynamical filtering. To test more systematically this
 296 relationship, we also calculated the correlations between measured and estimated MFs
 297 and for each flight (Table ??). To test the significance, we measure the number of De-
 298 grees of Freedom (DoF) present in each dataset, and calculate for that the decorrelation
 299 time scale, which we take as the lag in day beyond which the lag-autocorrelation of the
 300 series falls below 0.2. As this time-lag varies from one series to the other, we give explic-
 301 itly in column 5, the number of DoF, which is the duration of the flight divided by the
 302 decorrelation time scale. Note that for their decorrelation time, we consider for simplic-
 303 ity that evaluated with daily averaged observations, but found that it is not much dif-
 304 ferent from that evaluated with the offline estimates (not shown). In each case, we find

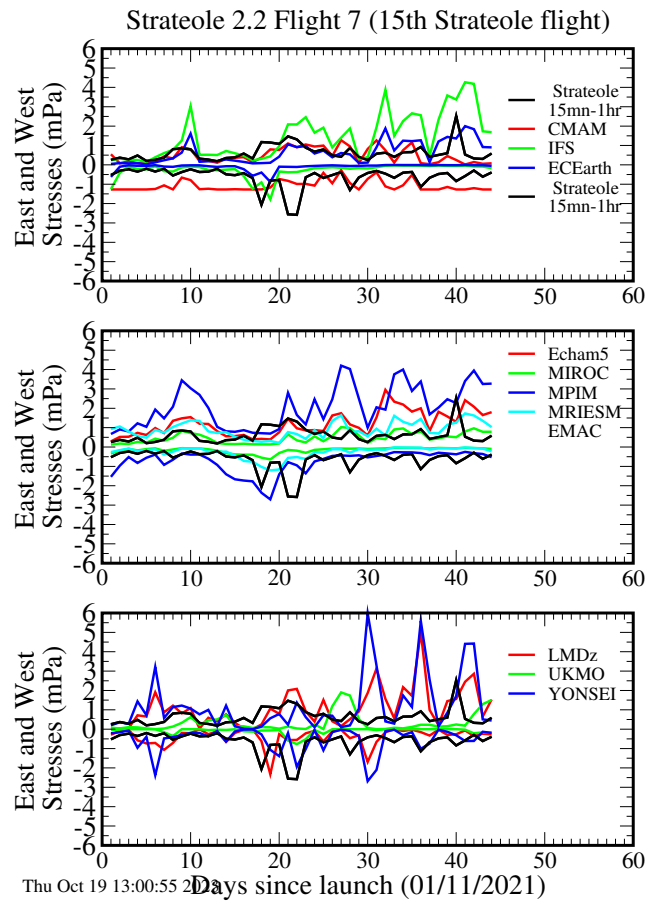


Figure 4. a) Same as Fig 3 but for Strateole 2 Phase 2 Flight 7.

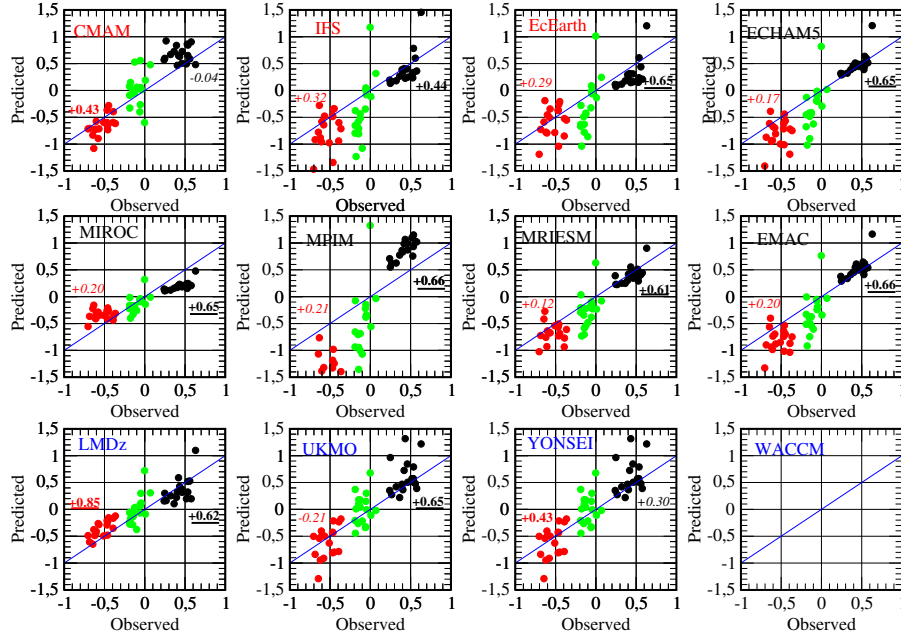


Figure 5. Scatter plot of the momentum fluxes measured by the balloon versus parameterized using different models. Only considered here the 18 balloon flights that last more than a month (East: black; West: red; Cumulated: green). Also shown are the correlations between observations and predictions, 99% significant levels are bold underlined, 95% are bold. Non significant values are shown in italic for information. The number of DoF for Pearson test is 18, about the number of balloon flights that last more than a month.

305 positive correlations, they are often significant in the Eastward direction the estimated
 306 westward being much more difficult to predict. In the eastward directions, differences
 307 seems more related to the launching altitude than to the nature of the scheme, CMAM
 308 launching waves near the tropopause whereas YONSEI often does so when clouds extent
 309 over the entire troposphere, in which case little dynamical filtering can occur between
 310 the launching level and the quite nearby flights.

311 A major differences between the scheme, is that in some the GW activity is related
 312 to convective heating or precipitation, whereas others consider uniform sources. In the
 313 first case its means that momentum fluxes could be underestimated in many circumstances,
 314 despite the fact that the amplitudes are realistic when considering long term averages.
 315 To analyse better this difference and its potential consequences, the Figure ?? presents
 316 PDFs of the distributions of the momentum fluxes considering all the daily data. For
 317 the PDFs (solid line), one sees that the balloons almost systematically measure fluxes
 318 with amplitude between 0.1mPa and 10mPa (see Figure ??a), whereas in the parame-
 319 terizations with convective/precipitation sources there are many more contributions from
 320 the smaller amplitude momentum fluxes (solid red), not mentioning that the zero val-
 321 ues are excluded from PDFs when plotted versus the logarithm of MF amplitudes.

322 4 Conclusion

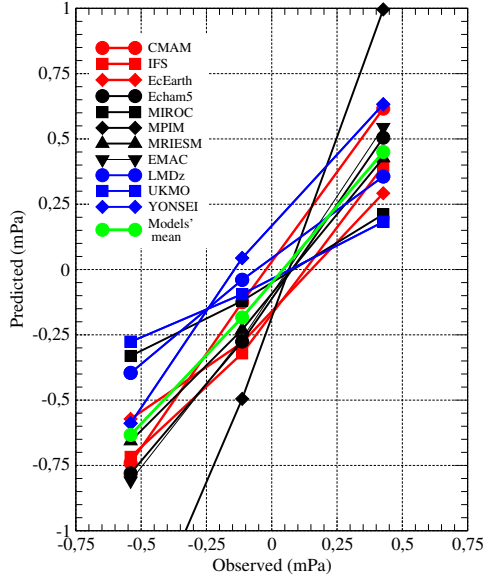


Figure 6. East, West and cumulated zonal momentum fluxes averaged over the Strateole 2 phase 1 and 2 period period and according to participating models.

East	Day Dof	CM AM	IFS	ECE ARTH	Ech am5	MI ROC	MPI M	MRI ESM	EM AC	LMD z	UK MO	YON SEI
Phase 1	670-216	-0.07	0.53	0.52	0.52	0.48	0.49	0.44	0.48	0.49	0.34	0.32
Phase 2	621-322	-0.19	0.41	0.38	0.38	0.33	0.34	0.30	0.33	0.40	0.34	0.20
Phase 1+2	1291-538	-0.11	0.49	0.47	0.45	0.41	0.41	0.36	0.40	0.46	0.34	0.27
West	Day Dof	CM AM	IFS	ECE ARTH	Ech am5	MI ROC	MPI M	MRI ESM	EM AC	LMD z	UK MO	YON SEI
Phase 1	670-216	0.14	-0.07	-0.07	-0.13	-0.03	-0.04	-0.04	-0.04	0.30	-0.03	0.10
Phase 2	621-322	0.21	0.18	0.16	0.03	0.00	0.01	0.05	-0.01	0.40	0.04	0.13
Phase 1+2	1291-538	0.17	0.05	0.04	-0.05	-0.02	-0.02	0.01	-0.02	0.34	0.00	0.11

Table 4. Correlation between observed and measured fluxes, strateole phases 1 and 2.

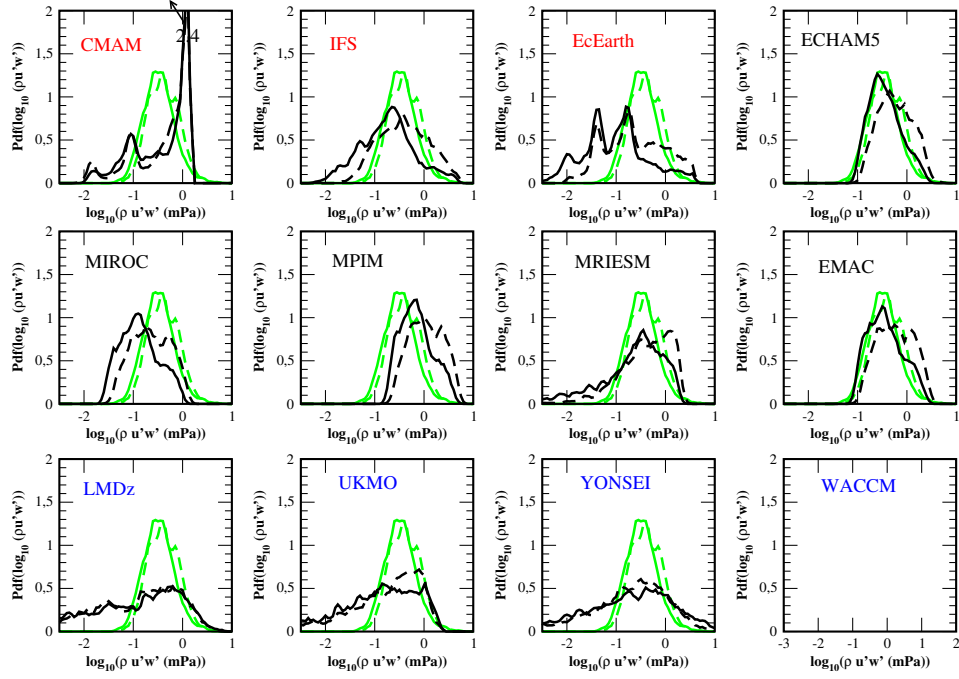


Figure 7. PDFs of daily values of Momentum flux distribution. The PDFs are calculated from histograms of 1291 MFs daily value within intervals of $\Delta(\log_{10} \overline{\rho u'w'}) = 0.05$, thereafter smoothed by a 5 point non-recursive filter with weight (0.1, 0.2, 0.4, 0.2, 0.1). Measured values are in green, estimations using ERA5 data and the parameterizations are in black. Solid lines are for Eastward, dashed lines are for Westward.

5 Open Research

Balloon data presented in (Haase et al., 2018) can be extracted from the STRATEOLE 2 dedicated web site: <https://webstr2.ipsl.polytechnique.fr>

ERA5 reanalysis data are described in (Hersbach et al., 2020) and can be extracted from the COPERNICUS access hub: <https://scihub.copernicus.eu/>

The LMDz-6A GCM used for CMIP6 project is described in (Hourdin et al., 2020), it can be directly installed from the dedicated webpage: <https://lmdz.lmd.jussieu.fr/utilisateurs/installation-lmdz>

Acknowledgments

This work was supported by the VESRI Schmidt Future project "DataWave".

Appendix: Running the offline code

To run the models parameterizations in offline mode and compare with daily values of momentum fluxes measured during strateole 2, download the file `offline_v9_Strateole_QBOi_Open.tar`, on the web page:

```
wget https://web.lmd.jussieu.fr/~flott/DATA/offline_v9_Strateole_QBOi_Open.tar.gz
```

Then gunzip and do `tar -xvf offline_v9_Strateole_QBOi_Open.tar`

In the directory, `offline_v9_Strateole_QBOi_Open`:

Run directory It basically contains a script that compile the programs, link to the input dataset and produce various outputs. The Makefile certainly needs to be adapted to the computer.

To launch predictions for Strateole-2 phase 1, launch: `./laun_ph1ball_gwd_era5.sh`
For phase 2, `ph1` → `ph2`.

Fortran Codes: all the fortran routines are located in `prog`.

laun_gwd_era5.f90: Main program loading input data in netcdf format and calculating drag and momentum fluxes at the balloon place.

preci_gwd_LMDz_QBOi.f90: LMDz Multiwaves routines predicting gwdrag from precipitation

gwsat_Modnam.f90: the globally spectral scheme using the (Warner & McIntyre, 1996)'s scheme version by J. Scinocca.

hinesgw6g_plus_subs.f HDS scheme

gw_ussp_core.f90: The WMI scheme with amplitude keyed to precipitation used in some UKMO runs.

cgwcalc.f90: Multiwave scheme developed at Yonsei's university

Input Data: All the input data are located in the directories `hourly_ph1` and `hourly_ph2` for phase 1 and 2 respectively. For instance, 1hr average of the strateole2 momentum fluxes are in

`ALL_STRATEOLE2_Balloon_ph1_1day15min.nc`
and

`ALL_STRATEOLE2_Balloon_ph1_1hrs15min.nc`

for the waves with periods between 1day and 15mn and between 1Hr and 15 mn respectively.

Still in this directory, the ERA5 reanalysis products, which include winds temperature, cloud liquid water, and surface log pressure, over a $5^\circ \times 5^\circ$ domain centered at the balloons drifting locations are in `Input_ERA5_data_all_variables_balloons_ph1.nc`. Precipitation every hours are also included. The diabatic heatings are from fore-

369 cast. All datas that are only provided every 3hr are linearly interpolated in time
 370 to give hourly values.

371 Output data (Part 1)

372 All the outputs are in the output_ph1 and output_ph2 directories:

373 **Netcdf:** contains the output of the schemes in netcdf format on the vertical col-
 374 umn and over the 5°x5° domain over which the ER5 data are provided. There is
 375 one netcdf dataset by balloons flight each contains output from all the schemes.

376 **Balloon_alt** After post processing by the python scripts launch_script_obs.py, are
 377 extracted the MFs at balloon flight altitude.

378 Python Scripts

379 A serie of Python scripts, located into python_script are proposed to compare the
 380 outputs of the scheme to the balloon data.

381 **launch_script_obs.py:** Reads the balloon flight data of MFs and averaged over
 382 1day and writte them in text format (ending with '.dat') and stored in **output/Balloon_alt/obs_output**

383 **launch_prediction_eachB_ysei.py:** extract from the prediction the values of the
 384 MFs at the balloons place and altitude. Results stored in text format (".dat" in
 385 **Balloon_alt/Pred_output_Balloon_altitude/**.

386 The next python scripts are cosmetic in the sense that they use the above two datasets
 387 to make plots of timeseries balloon averaged values, evaluate correlations, and his-
 388 tograms.

389 **timeseries_obs_pred_plot_all.py** Produces a lot of time series for each model
 390 and flights.

391 Output data (Part2) As a result, you can visualize timeseries of each flight here:

392 **output_ph1/Balloon_alt/figure_timeseries**

393 Histograms here: **output_ph1/histo**

394 Scatter plots and correlations here **output_ph1/correlation**

395 For phase 2, change ph1 in ph2.

396 xmgrace Alternative to calculate these diagnostics using fortran programs and xmgrace,
 397 the programs permit to combine statistics over the 2 phases of Strateole2.

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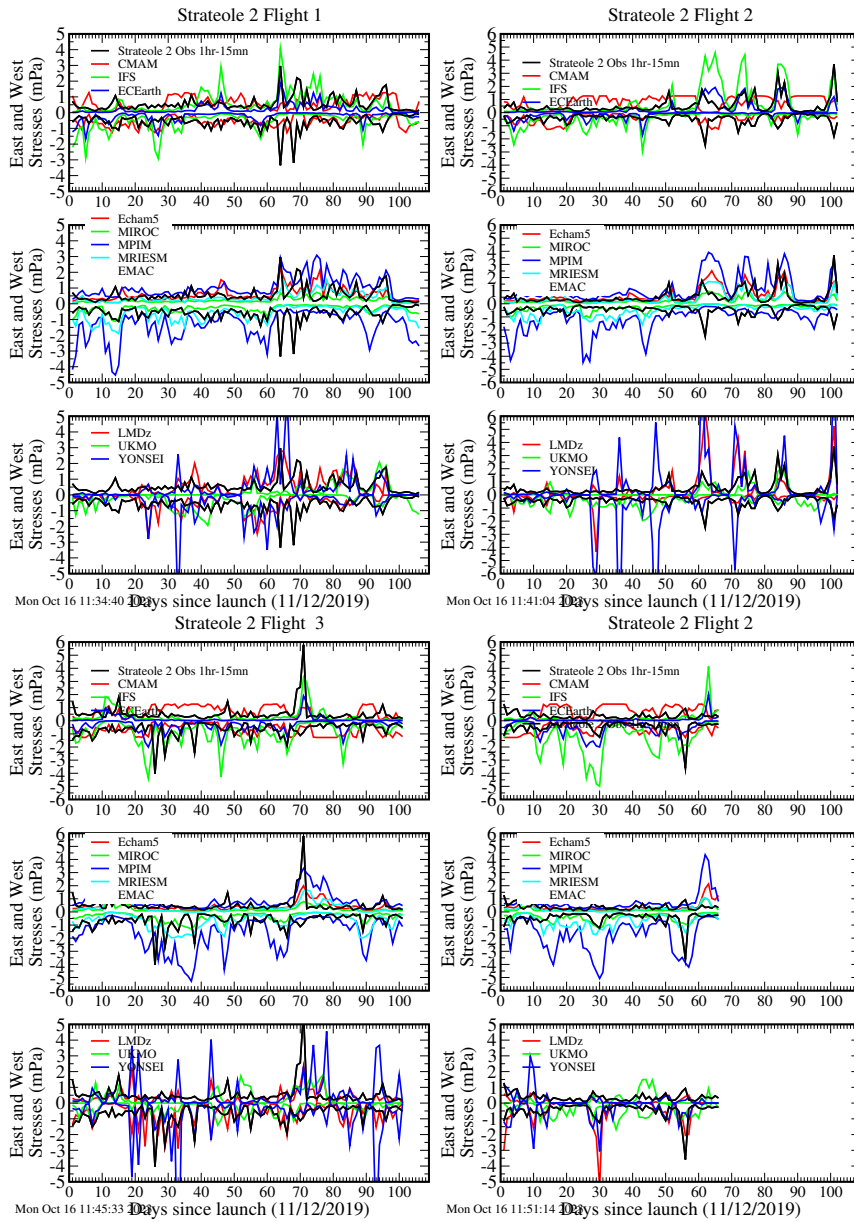
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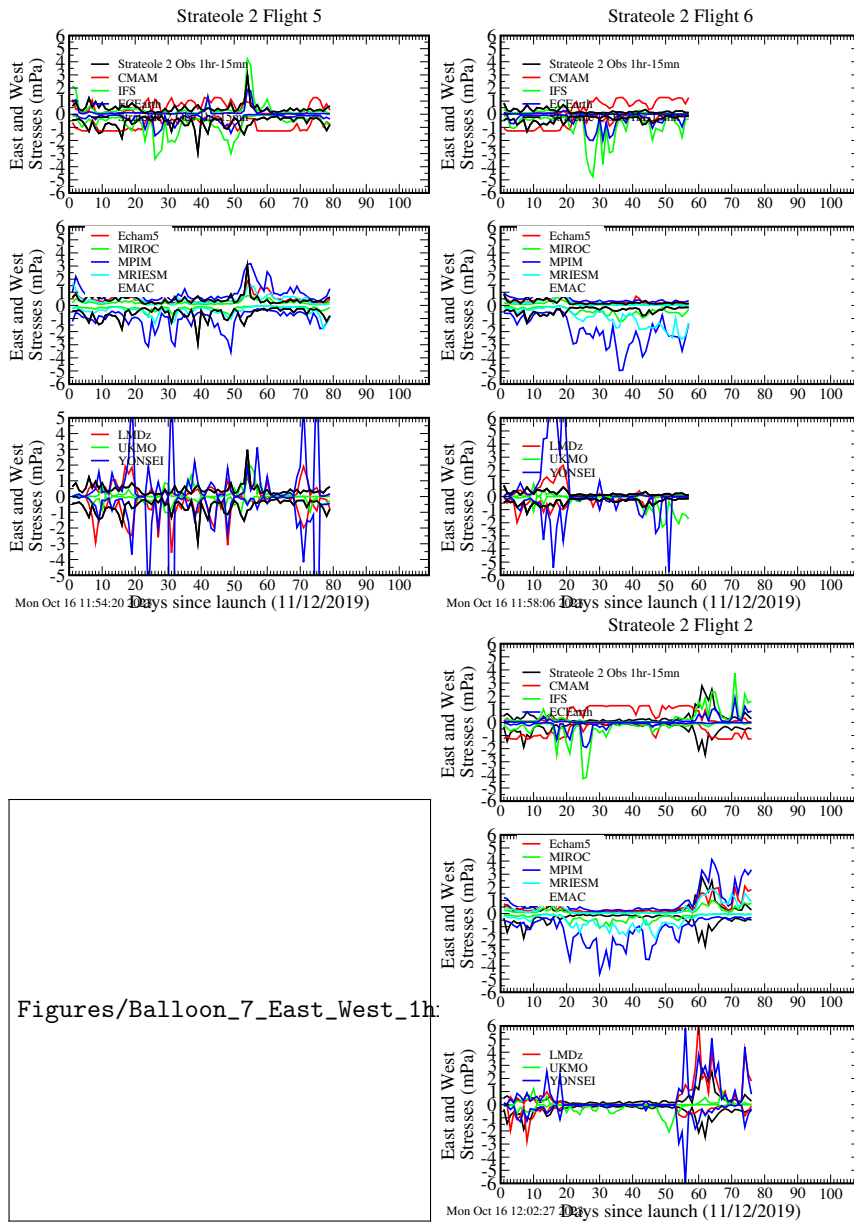
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Supplementary: Phase 1 flights and models

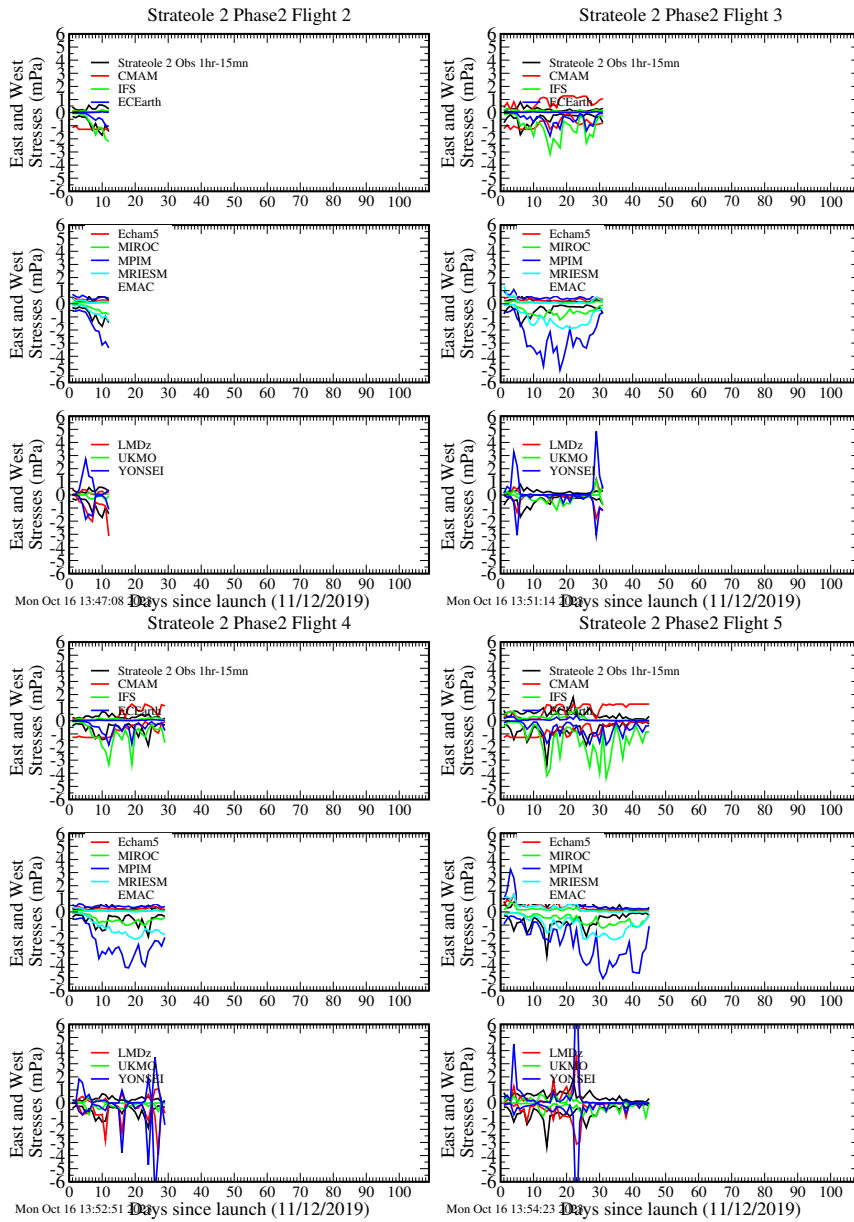


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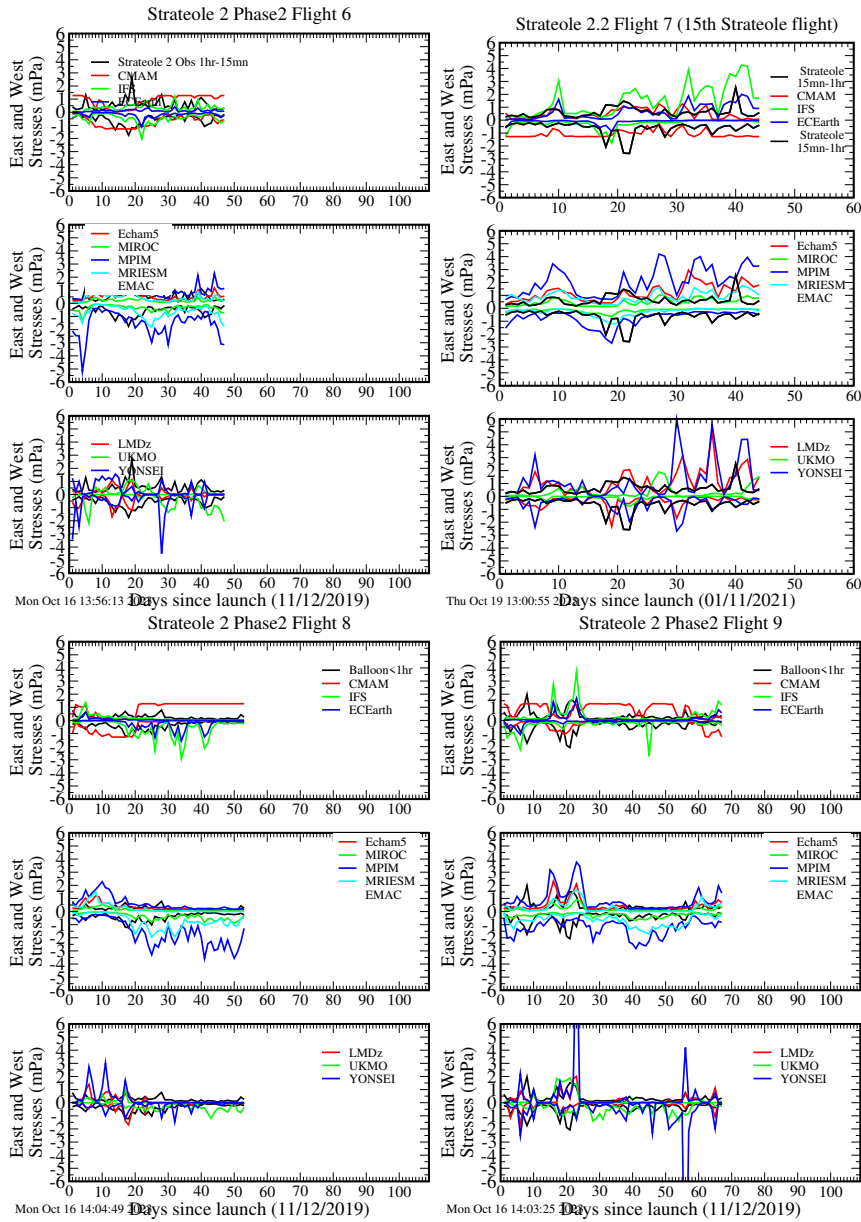
Supplementary: Phase 1 flights and models (continued)



Supplementary: Phase 2 flights and models

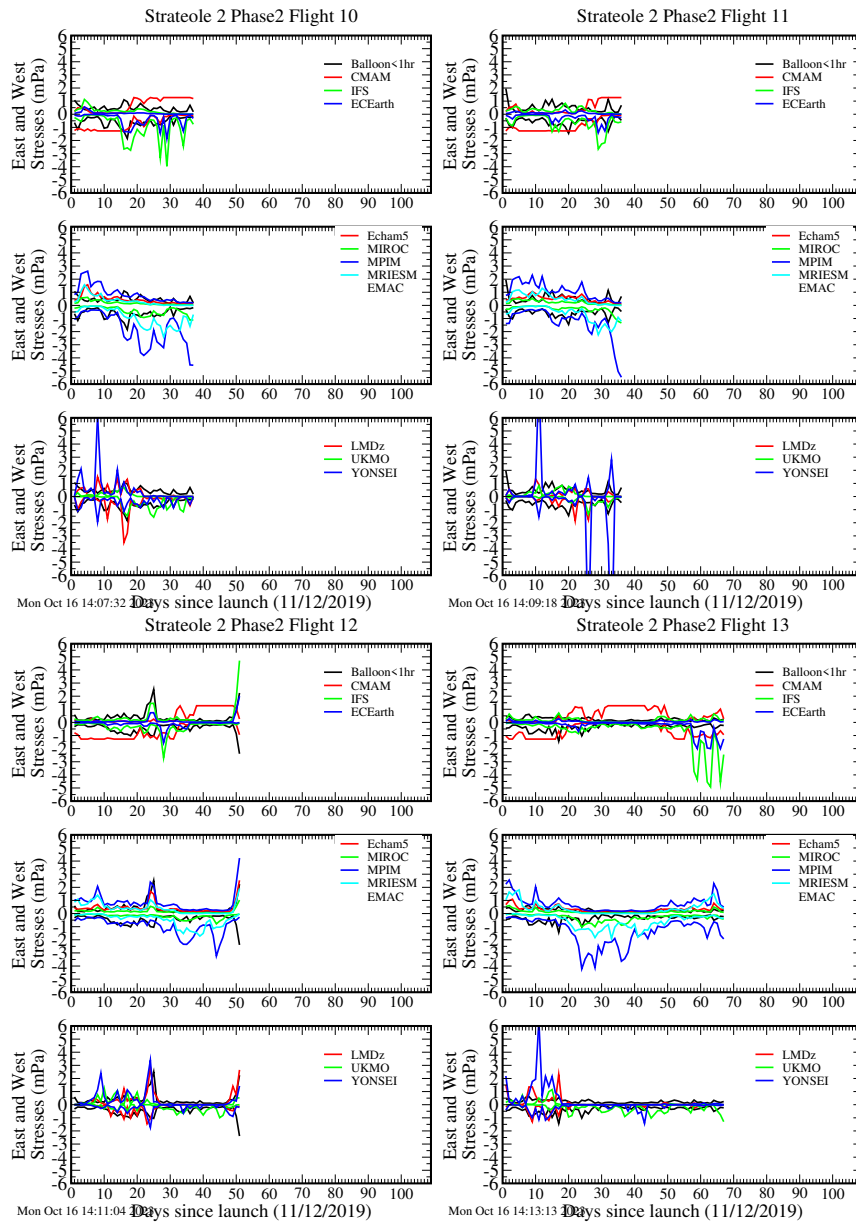


Supplementary: Phase 2 flights and models (continued)



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Supplementary: Phase 2 flights and models (continued)



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Supplementary: Phase 2 flights and models (continued)

