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EXOTOP December 4th and 5th at LATMOS, Guyancourt, France

Turbulence and Drags from Non-orographic Gravity Waves in the Mars Planetary Climate Model

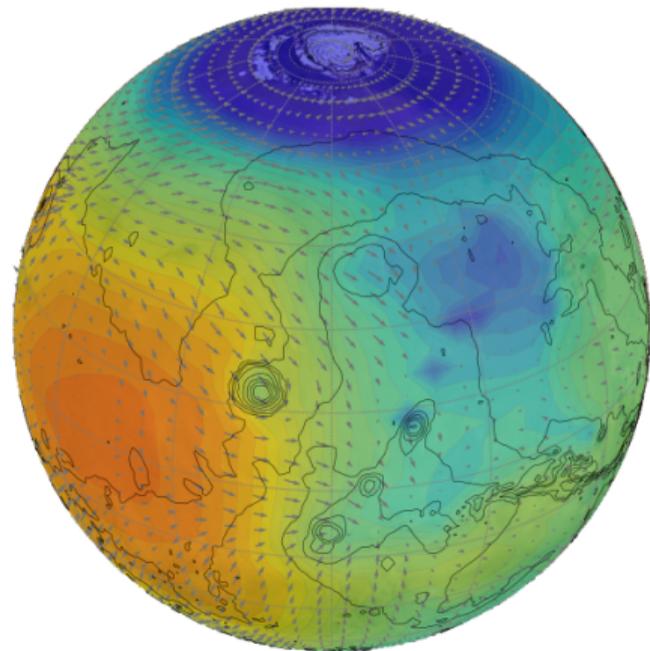
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Context

- Introduction
- Formalism (Drags+Turbulence)
- Simulations v.s. Observations



small waves, big effects

Non-orographic GWs:

- > Waves EP-flux
- >> Divergences of EP-fluxes (Drags)
- >>> Induced turbulence/mixing
- >>>> Mixing fluxes induced drags

Planetary waves

- > β effect, $\mathfrak{S}(m)$, fast decay

Thermal tides

- > Forced planetary-scale GWs

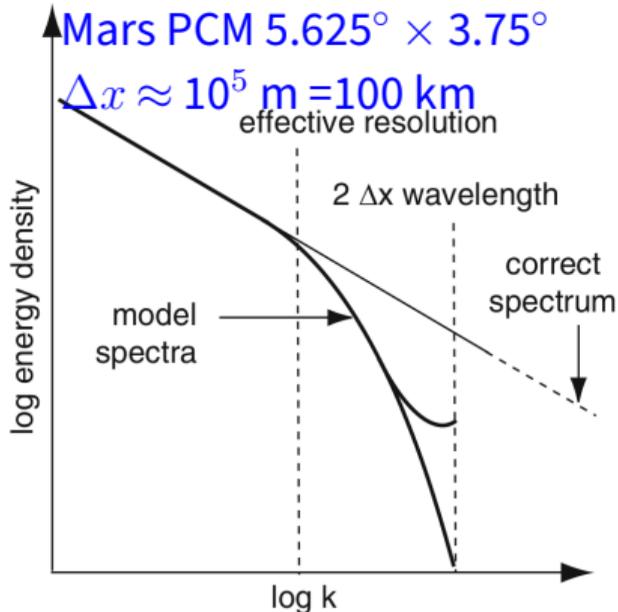
Orographic GWs:(bugged)

- > Drags from Blocked Flows
- > Propagating Waves from Top of Blocked Flows
- >> induced Mixing (not parameterized, see tomorrow's talk)

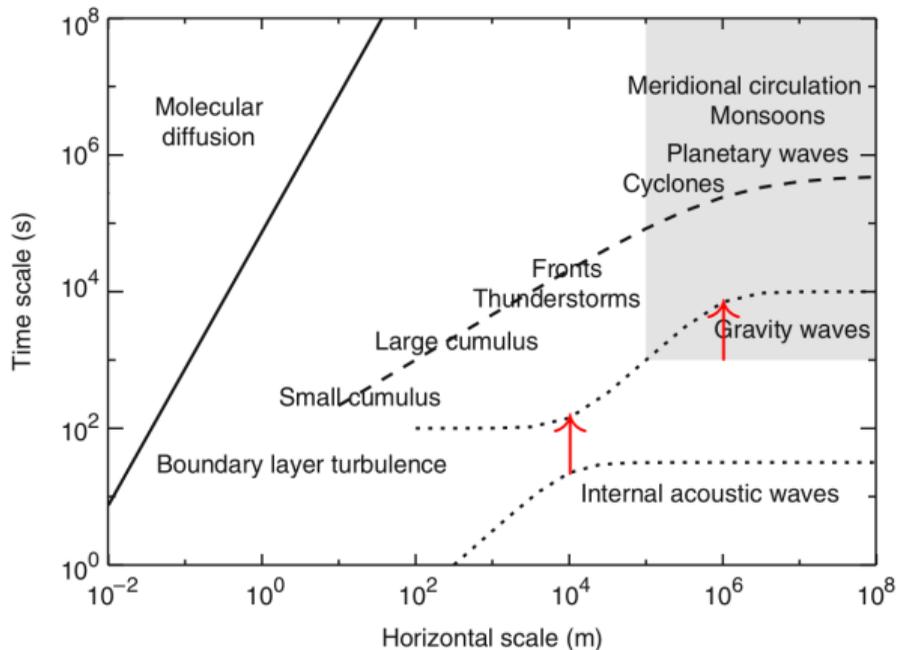
No inertial filtering is considered currently, thus these schemes do not apply to fast-rotated Planets such as Jupiter, Saturn, etc; See discussion demain

Gravity Waves in Mars PCM

$$k_{eff} \in [2\pi/10\Delta x, 2\pi/4\Delta x]$$



Model physical $\delta t \approx 4.6 \times 10^2 \text{ s}$



Lauritzen et al.(2011)

Modeled gravity waves: 10^4 - 10^6 m
 with life cycle of 88775 s

Setting Up

Variables

Ranges

wavenumber k, l [m^{-1}]

$$2\pi/10\Delta x \leq k \leq \max(\pi/\sqrt{\Delta x \Delta y}, 2\pi/\Delta x/2)$$

phase velocity c [m s^{-1}]

$$c = |\vec{u}_r|, 2|\vec{u}_r|$$

saturation factor S_c

$$1 < S_c \leq 3 \text{ (due to mixing)}$$

viscosity-eddy μ^*

$$\mu \bar{D}_{eddy}(z_b)$$

source altitude z_r

ave(PBL), vPBL, other

max EP-flux E_{max}

$$2\Omega \sin \phi \times \frac{\overline{v'\theta'}}{\theta_z} \text{ (RMC by GWs at } z_b)$$

wavelength: $\lambda = 2\pi/k$; frequency: $\omega = kc$

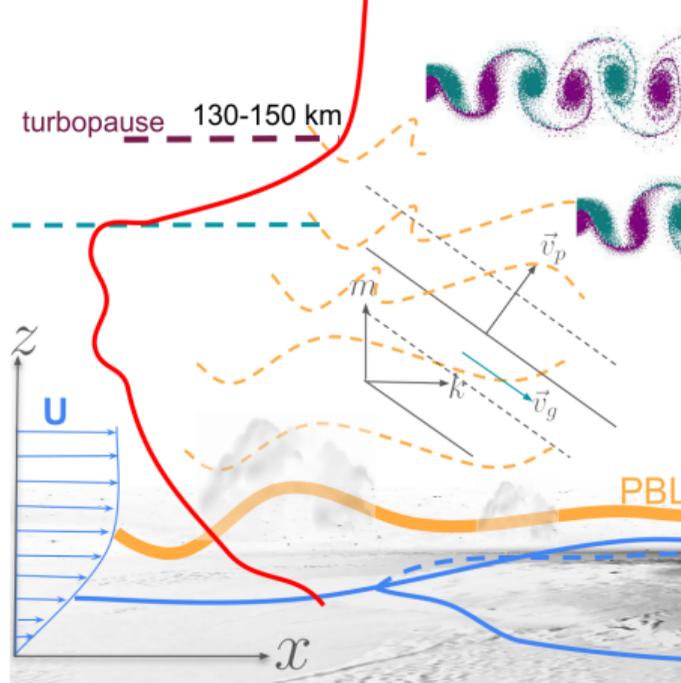
EP-Flux

Pseudo-Momentum:

$$\vec{E}_j^{zr}(k_j, l_j, \omega_j) = \Re \left\{ \rho_r \frac{\hat{u} \hat{w}^*}{2} \right\}$$

$$\hat{u} = -\frac{m}{|\vec{k}|} \hat{w}$$

$$\vec{E}_j^{zr}(k_j, l_j, \omega_j) = -\rho_r \frac{\vec{k}}{2|\vec{k}|} m_j(z_r) |\hat{w}_j(z_r)|^2 \frac{\vec{k} \Omega}{|\vec{k} \Omega|}$$



turbulence

$$\hat{w}_j(z_{l+1}) = \underbrace{\Theta[\Omega(z_{l+1}) \times \Omega(z_l)]}_{\text{critical-layer}} \text{Min} \left\{ \hat{w}_j(z_l) \sqrt{\frac{m_r(z_l)}{m_r(z_{l+1})}} e^{-\frac{\mu}{\rho} \left(\int_{z_l}^{z_{l+1}} m_i^{ave} d\zeta \right)}, \underbrace{\hat{w}_{j,s}}_{\text{saturation}} \right\}$$

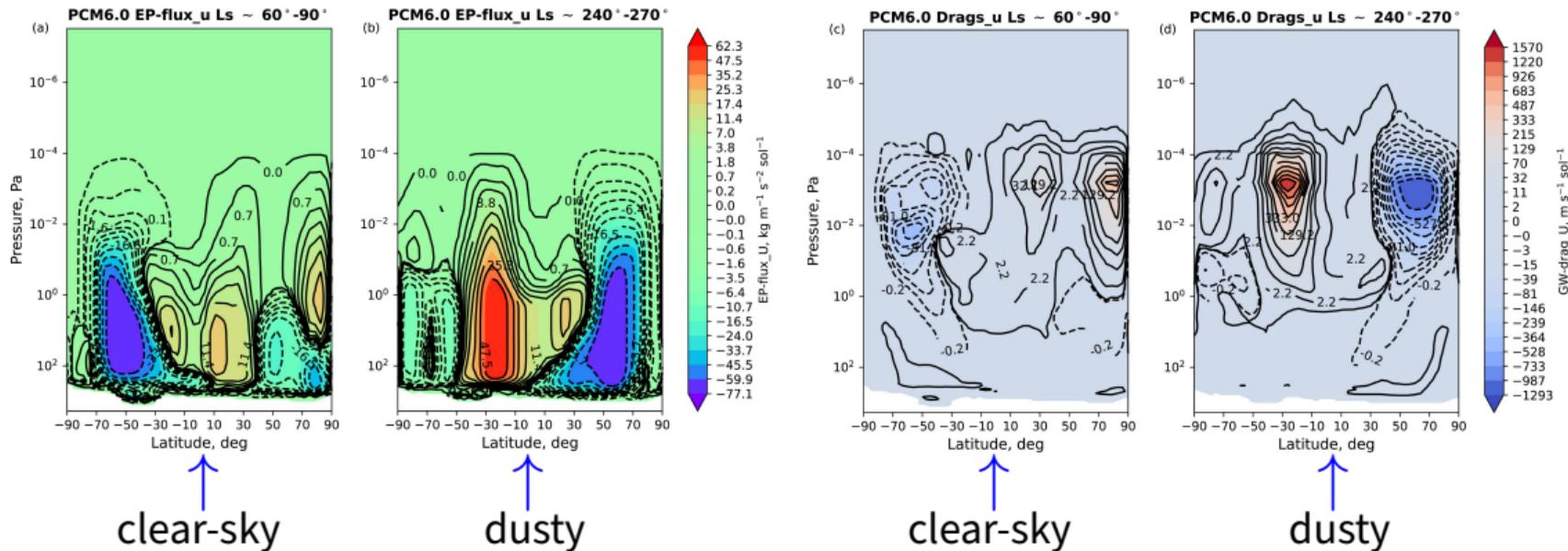
↑
kinematic viscosity

Lott et al.(2012),2013

How Do Non-orographic Gravity Waves Propagate ?

EP-flux [$\text{kg m}^{-1} \text{s}^{-2} \text{sol}^{-1}$]

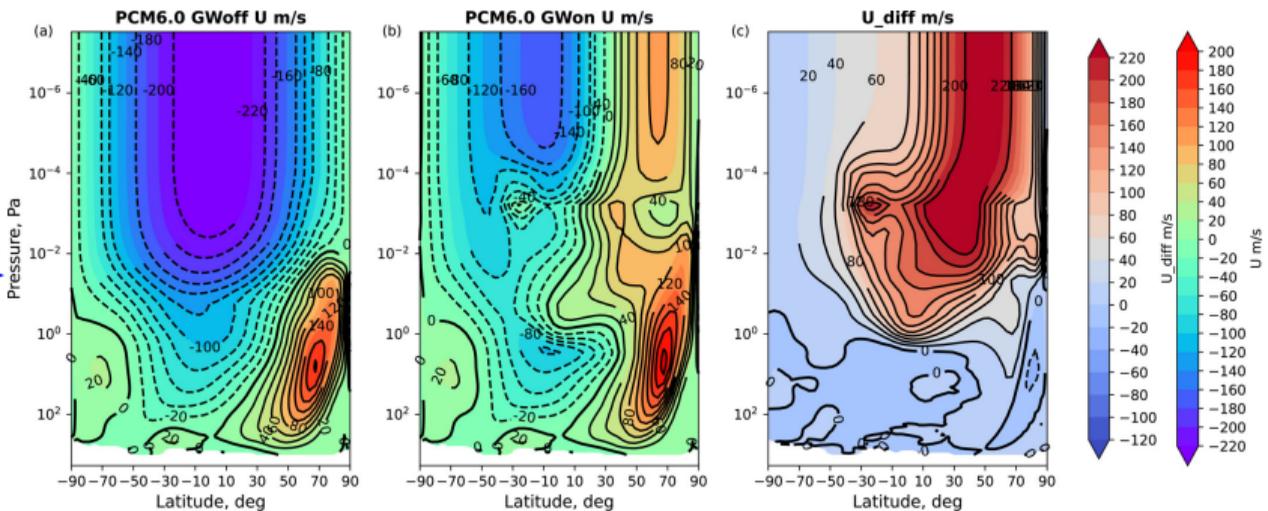
Drags [$\text{m s}^{-1} \text{sol}^{-1}$]



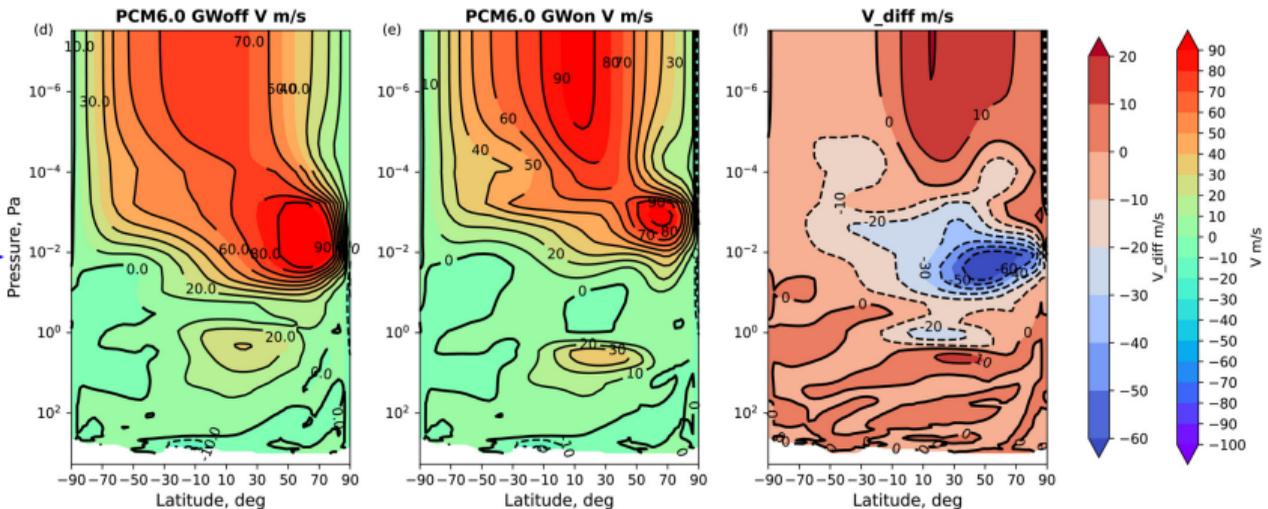
Liu et al.(2023),JGR-Planets:10.1029/2023JE007769

Effects on Flows

zonal winds \Rightarrow



meridional winds \Rightarrow



Drag on Circulation

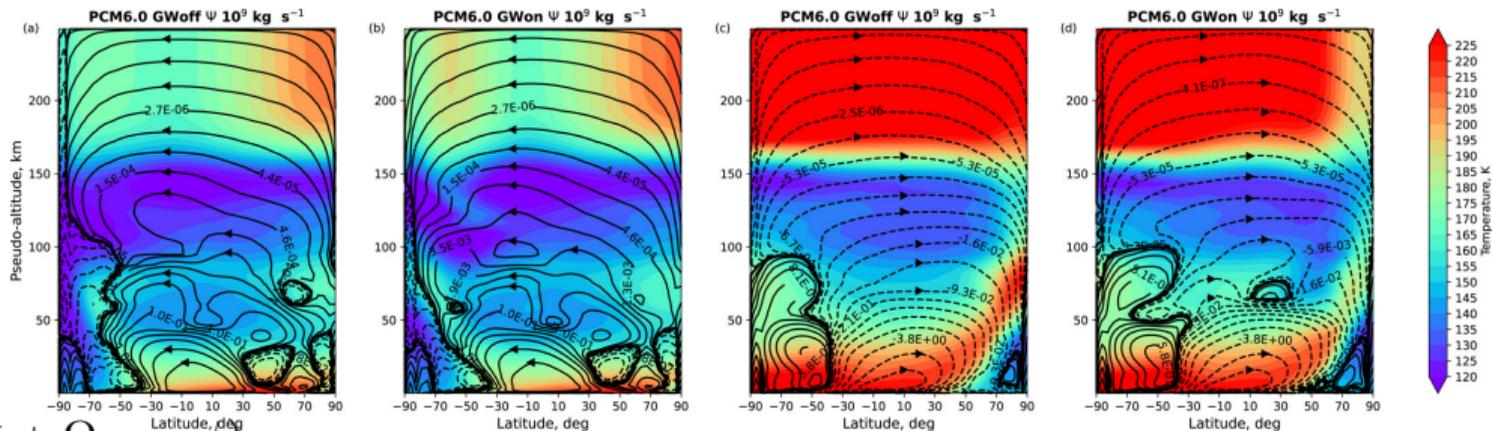
$$\Psi = \frac{2\pi a \cos \phi}{g} \int_0^p \bar{v} dp'$$

clear-sky

dusty

$$\Phi$$

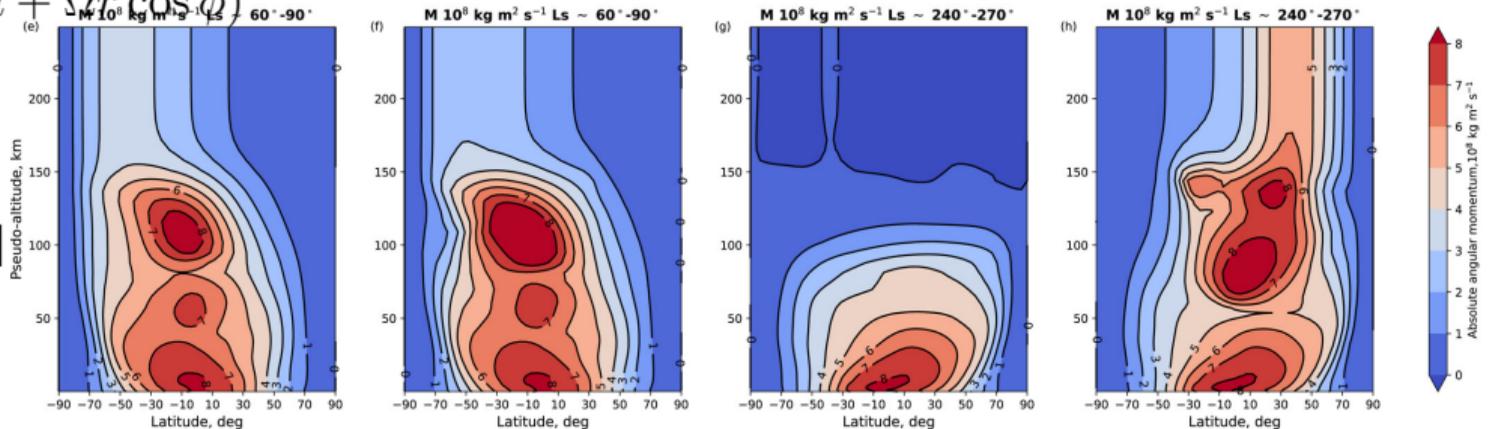
$[10^9 \text{ kg s}^{-1}]$



$$M = r \cos \phi (\bar{u} + \Omega r \cos \phi)$$

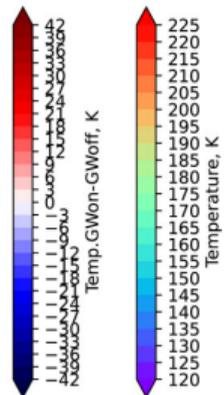
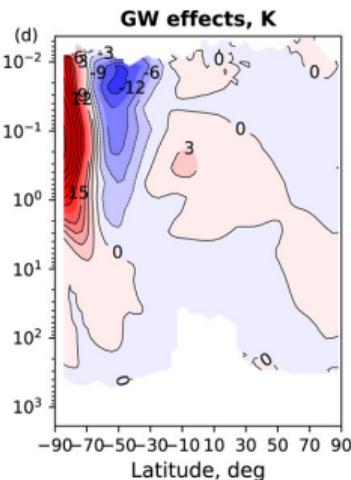
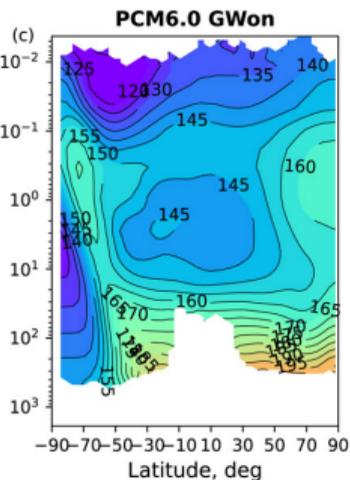
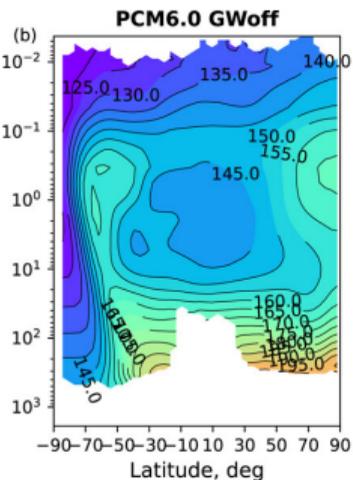
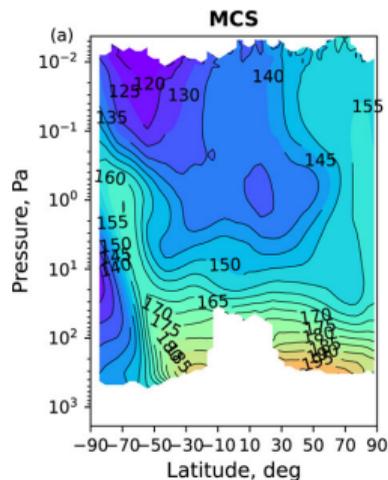
$$M$$

$[10^9 \text{ kg m}^2 \text{ s}^{-1}]$

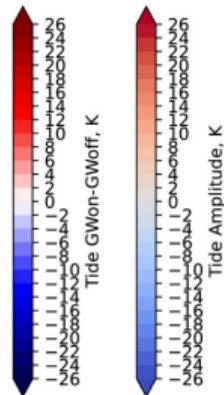
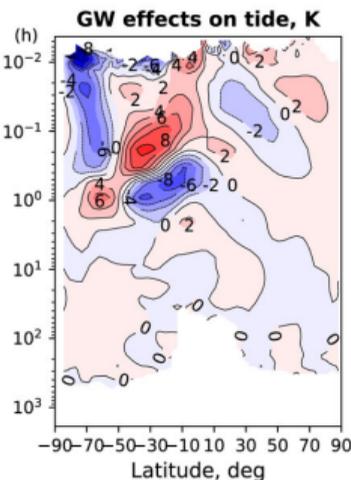
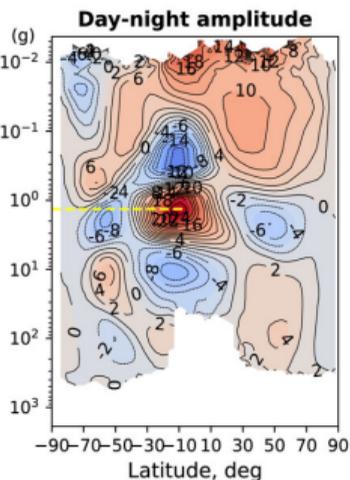
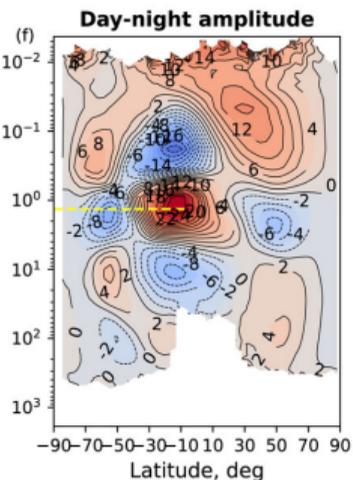
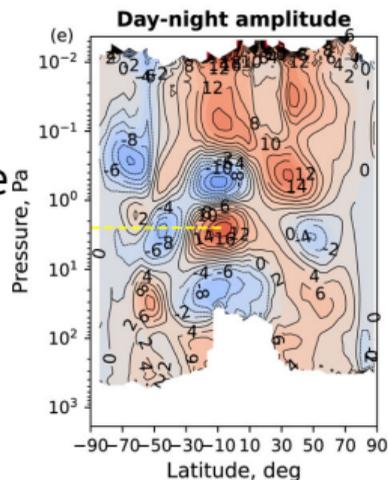


clear-sky

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[K]

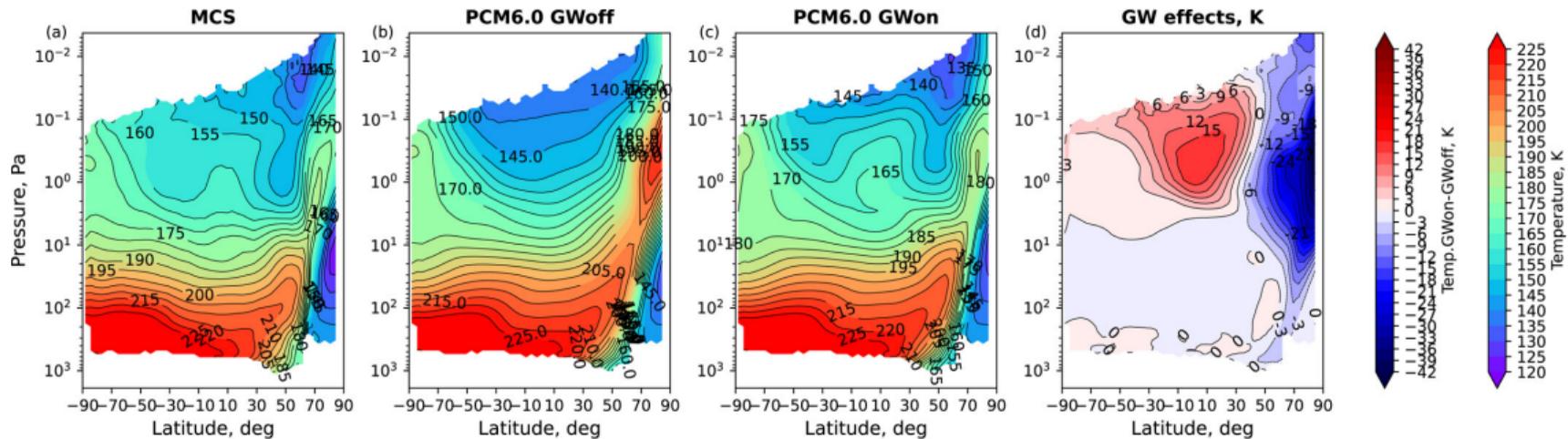


Tide
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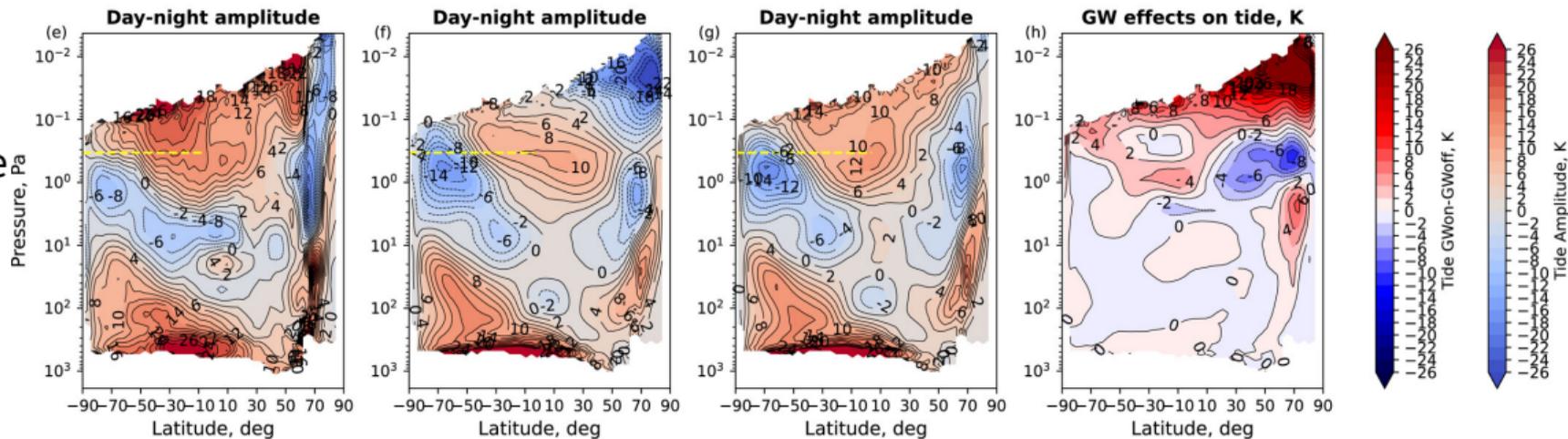


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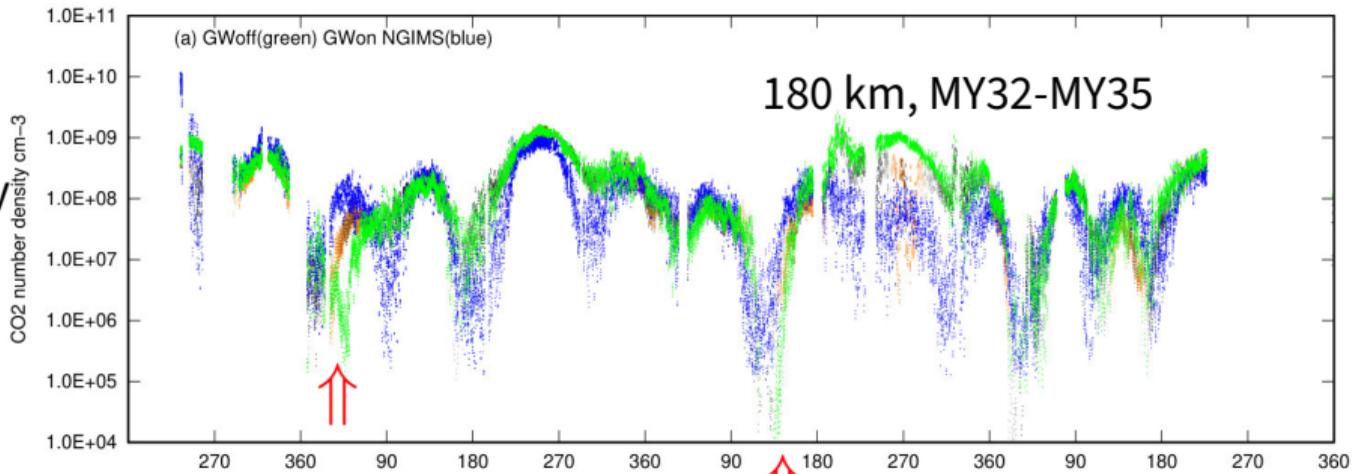


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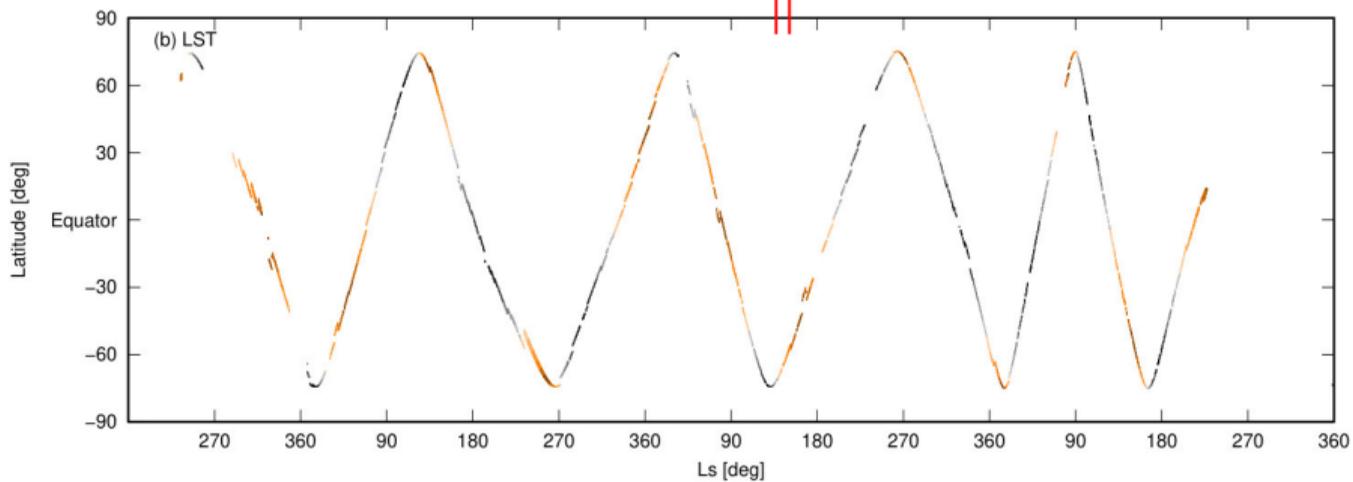


NGIMS/GWoff/GWon

CO2 number density



Local Solar Time:
day/night



Turbulence from Waves

Assumption 1: Turbulence triggered by the release of wave momentum, viz.

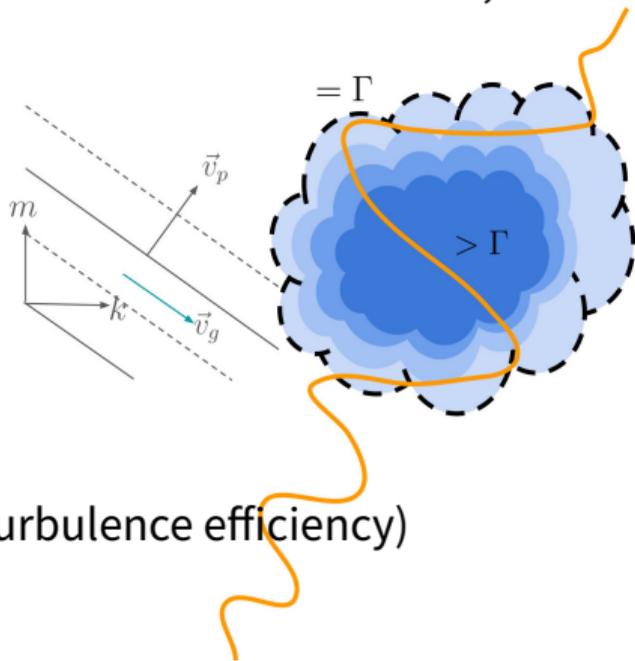
$$\text{Turbulence} \propto \frac{1}{\rho} \nabla(\vec{E})$$

Assumption 2: Non-superadiabatic Principle

$$\left| \frac{dT}{dz} \right|_{gw} = \Gamma \text{ at } z_b \text{ (Lindzen 1981)}$$

Assumption 3: Symmetric Principle

Assumption 4: Linear Damping (instability-to-turbulence efficiency)



Wave Thermodynamics: m_i at z_b

Assumption 2: Non-superadiabatic Principle

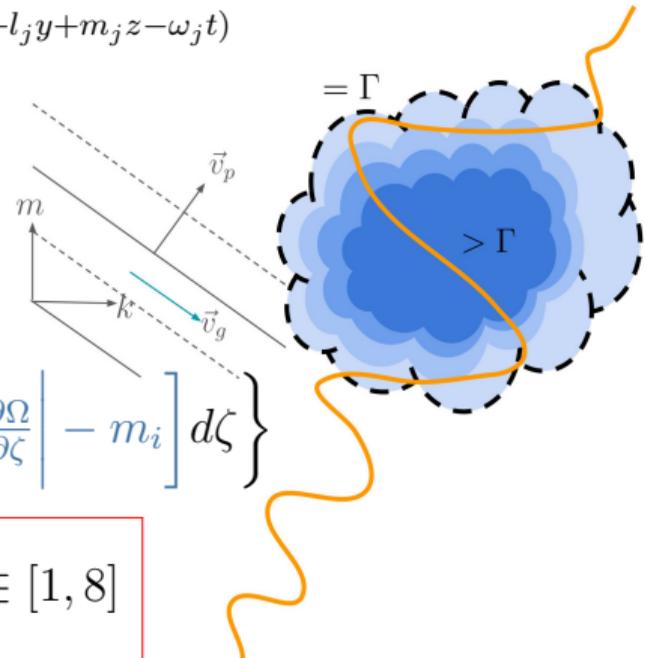
$$-i\Omega T' + w'\Gamma = 0 \quad w'_j = \hat{w}_j(z) \underbrace{e^{z/2H}}_{\rho(z)} e^{i(k_j x + l_j y + m_j z - \omega_j t)}$$

$$i\vec{k}_j(c - \bar{u})\delta T = \hat{w}_j e^{z/2H} \Gamma$$

$$\Re\left\{\frac{d\delta T}{dz}\right\} = -\frac{\Gamma A |m_r|^{3/2} e^{\int_0^z (\frac{1}{2H} - m_i) d\zeta}}{N|\vec{k}|}$$

$$\Re\left\{\frac{d\delta T}{dz}\right\} = -\frac{\Gamma A}{N|\vec{k}|} \exp\left\{\int_0^z \left[\frac{1}{2H} + \frac{3}{2} \left|\frac{1}{N} \frac{\partial N}{\partial \zeta} - \frac{1}{\Omega} \frac{\partial \Omega}{\partial \zeta}\right| - m_i\right] d\zeta\right\}$$

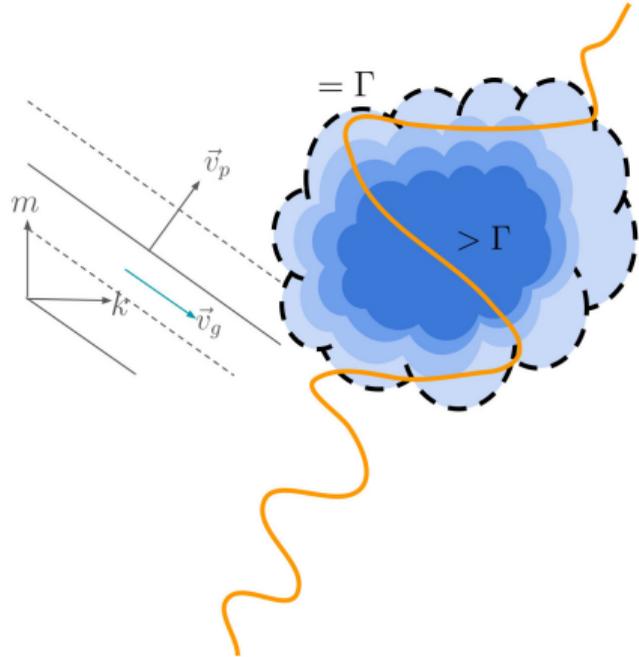
$$\frac{1}{2H} + \frac{3}{2} \left|\frac{1}{N} \frac{\partial N}{\partial z} - \frac{1}{\Omega} \frac{\partial \Omega}{\partial z}\right| = m_{i,s}, \quad z = z_b, \forall j \in [1, 8]$$



Wave Thermodynamics: NSP

Assumption 2: NSP $\left| \frac{dT}{dz} \right|_{gw}^2 = \Gamma^2$

$$\Re \left\{ \frac{d\delta T}{dz} \right\} = - \frac{\Gamma A |m_r|^{3/2} e^{\int_0^z (\frac{1}{2H} - m_i) d\zeta}}{N |\vec{k}|}$$



Saturation altitude z_b and A^2

$$A^2 m_r^{-1} \exp \left(- 2 \int_0^{z_b} m_i d\zeta \right) = m_r^{-4} N^2 |\vec{k}|^2 \exp \left(- \int_0^{z_b} \frac{1}{H} d\zeta \right), \quad z = z_b$$

Wave Thermodynamics: Saturated EP-flux

Assumption 2: Non-superadiabatic Principle $\left| \frac{dT}{dz} \right|_{gw}^2 = \Gamma^2$

$$A^2 m_r^{-1} \exp \left(-2 \int_0^{z_b} m_i d\zeta \right)$$

$$= m_r^{-4} N^2 |\vec{k}|^2 \exp \left(- \int_0^{z_b} \frac{1}{H} d\zeta \right) \implies |\hat{w}_j(z_b)|^2 \approx A^2 m_r^{-1} \exp \left(-2 \int_0^{z_b} m_i d\zeta \right)$$

$$z \in [z_b, z_b + \delta z_u]$$

$$= m_r^{-4} N^2 |\vec{k}|^2 \exp \left(- \int_0^{z_b} \frac{1}{H} d\zeta \right)$$

$$= \frac{\Omega^4}{N^2 |\vec{k}|^2} \exp \left(- \int_0^{z_b} \frac{1}{H} d\zeta \right)$$

$$\vec{E}_j^{zu+1} = \frac{\vec{k}\Omega}{|\vec{k}||\Omega|} \Theta[\Omega(z_{u+1})\Omega(z_u)]$$

$$\text{Min} \left\{ |\vec{E}_j^{zu}| e^{-2\frac{\mu m_i}{\rho} \delta z}, \rho_r S_c^2 e^{-\frac{z_{ave}}{H}} \frac{|\Omega|^3 k^{*2}}{2N |\vec{k}|^4} \right\}$$



$$\hat{w}_{j,s} = \frac{\Omega^2}{|\vec{k}|N} e^{-z/2H} S_c \frac{k^*}{|\vec{k}|}$$



Wave Thermodynamics: Linear Damping

Assumption 4: Linear Damping

$$D_{eddy}^j \frac{\partial}{\partial z^2} \begin{Bmatrix} \bar{u} \\ \delta T \end{Bmatrix} = -m_r^2 D_{thermal}^j \begin{Bmatrix} \bar{u} \\ \delta T \end{Bmatrix}$$

$$Prandtl = D_{eddy}^j / D_{thermal}^j = 1$$

$$[i\vec{k}(c - \bar{u}) + m_r^2 D_{eddy}] \delta T = w' \Gamma$$

$m = m_r + im_i$, we rewrite $c = c_r + ic_i$,

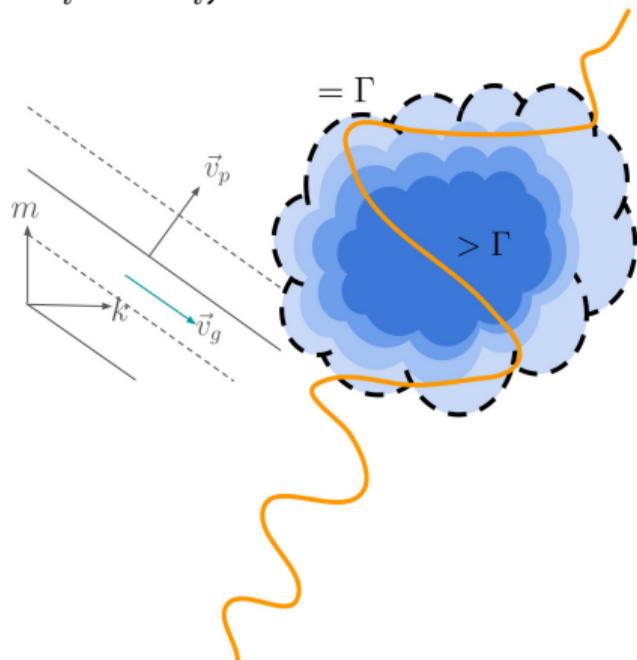
$$\vec{k}c_i = m_r^2 D_{eddy}^j$$

Wave Thermodynamics: Linear Damping

Assumption 4: Linear Damping: The turbulence m_i and c_i ,

$$\begin{aligned}
 m &= \frac{N|\vec{k}|}{\vec{k}(\bar{u} - c)} & c &= c_r + i c_i \\
 & & \leftarrow & \vec{k} c_i = m_r^2 D_{eddy}^j \\
 &= \frac{N|\vec{k}|}{\vec{k}(\bar{u} - c_r) - i m_r^2 D_{eddy}^j} \\
 &= -\frac{N|\vec{k}|\Omega}{\Omega^2 + \cancel{m_r^4 D_{eddy}^2}} + i \frac{N|\vec{k}| m_r^2 D_{eddy}^j}{\Omega^2 + \cancel{m_r^4 D_{eddy}^2}}
 \end{aligned}$$

$$m_i = \frac{N|\vec{k}| m_r^2 D_{eddy}^j}{\Omega^2} = \frac{m_r^3 D_{eddy}^j}{\Omega}, \quad z \in [surface, top]$$



Eddy Diffusion Coefficient $z = z_b$

Since

$$m_i = \frac{N|\vec{k}|m_r^2 D_{eddy}^j}{\Omega^2}, z \in [surf, top]$$

Saturated vertical wavenumber of turbulence,

$$\frac{1}{2H} + \frac{3}{2} \left| \frac{1}{N} \frac{\partial N}{\partial z} - \frac{1}{\Omega} \frac{\partial \Omega}{\partial z} \right| = m_{i,s}, \quad z = z_b$$

$$D_{eddy}^j = \frac{\Omega^4}{N^3|\vec{k}|^3} \left(\frac{3}{2} \left| \frac{1}{N} \frac{dN}{dz} - \frac{1}{\Omega} \frac{d\Omega}{dz} \right| + \frac{1}{2H} \right)$$

$$D_{eddy}^j = S_{mix} \frac{\Omega^4}{N^3|\vec{k}|^3} \left(\frac{3}{2} \left| \frac{1}{N} \frac{dN}{dz} - \frac{1}{\Omega} \frac{d\Omega}{dz} \right| + \frac{1}{2H} \right), \quad z = z_b$$

Eddy Diffusion Coefficient $z > z_b$

Assumption 1: Turbulence triggered by momentum releasing, viz,
Turbulence $\propto \nabla(\vec{E})$

$$\begin{aligned} -\frac{1}{\rho} \frac{\partial \vec{E}_j^z}{\partial z} &= \alpha_{eff} \frac{\Omega^3}{N|\vec{k}|^2} m_i \frac{\vec{k}}{|\vec{k}|} \\ &= \alpha_{eff} \frac{\Omega^3}{N|\vec{k}|^2} \left(\frac{3}{2} \left| \frac{1}{N} \frac{dN}{dz} - \frac{1}{\Omega} \frac{d\Omega}{dz} \right| + \frac{1}{2H} \right) \frac{\vec{k}}{|\vec{k}|}, \quad z > z_b \end{aligned}$$

$$D_{eddy}^j = \alpha_{eff} \frac{\Omega}{N^2|\vec{k}|} \left\{ -\frac{1}{\rho} \frac{\partial \vec{E}_j^z}{\partial z} \right\}, \quad z > z_b$$

Eddy Diffusion Coefficient $z < z_b$

Assumption 3: Symmetric Principle: empirical fit

$$D_{eddy}^j = \mathcal{B} \times n_q^{-1/2}$$

Since we know $n \propto \exp -\frac{z}{H}$,

$$D_{eddy}^j = D_{eddy}^j(z_b) e^{\beta_{diff}(z-z_b)/H}, \quad z < z_b$$

General Equation

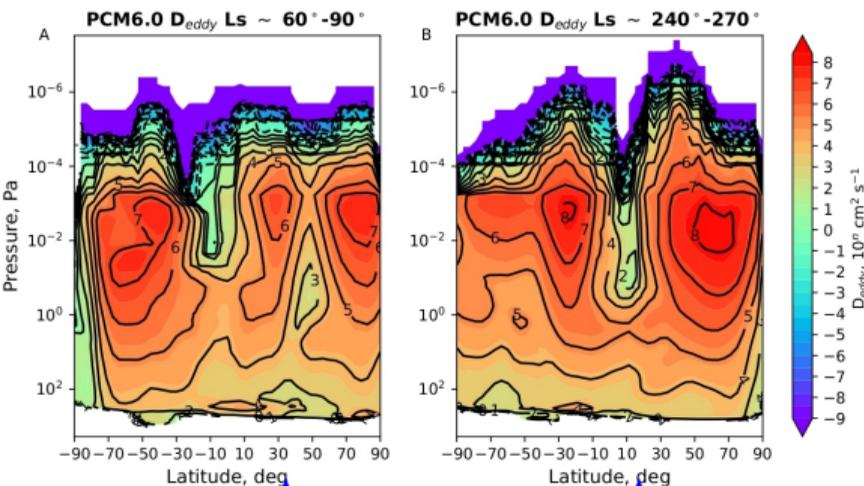
Eddy diffusion coefficient $\forall j$ th monochromatic wave:

$$D_{eddy}^j = \begin{cases} \text{Min} \left[\alpha_{eff} \frac{\Omega}{N^2 |\vec{k}|} \left(-\frac{1}{\rho} \frac{\partial \vec{E}}{\partial z} \right), S_{mix} \frac{\Omega^4}{N^3 |\vec{k}|^3} \left(\frac{3}{2} \left| \frac{1}{N} \frac{dN}{dz} - \frac{1}{\Omega} \frac{d\Omega}{dz} \right| + \frac{1}{2H} \right) \right] & , z \geq z_b \\ D_{eddy}(z_b) \exp \left[\beta_{diff} \frac{(z-z_b)}{H} \right] & , z < z_b \end{cases}$$

$$\alpha_{eff} = S_{mix} = 0.1 \text{ and } \beta_{diff} = \frac{3}{2}$$

Eddy Diffusion Coefficient

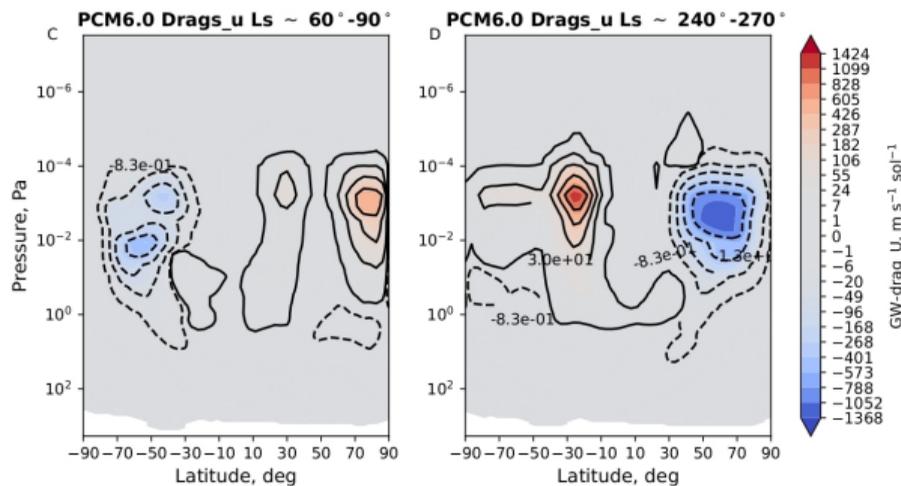
$$D_{eddy} = \sum_{j=1}^M D_{Eddy}^j, M = 8$$



↑
clear-sky

↑
dusty

Drags (GW drags not mixing drags)

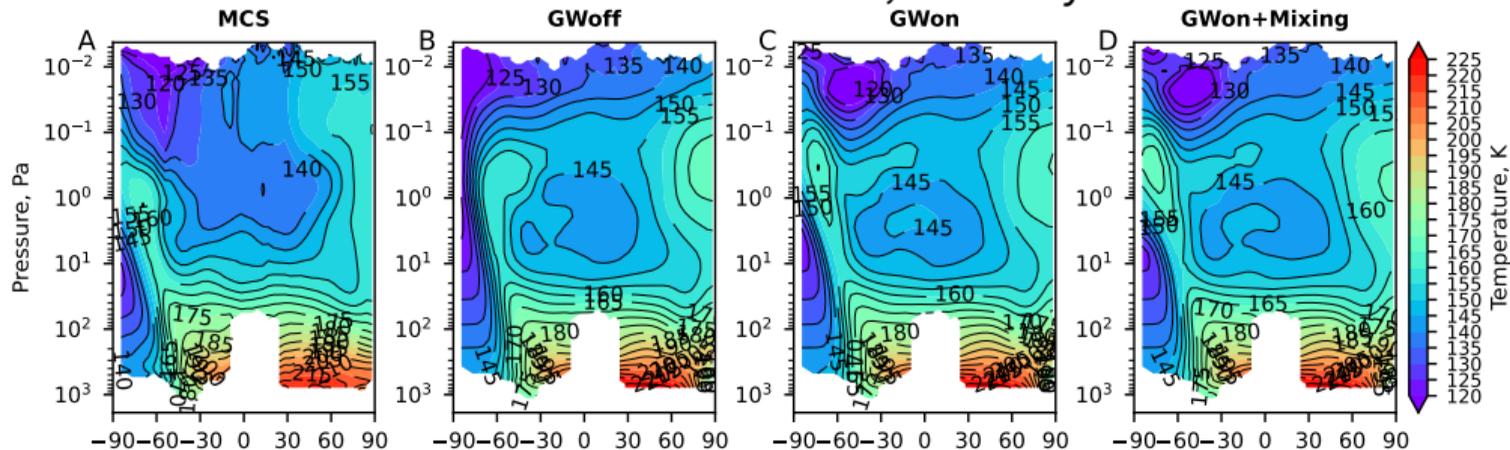


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clear-sky

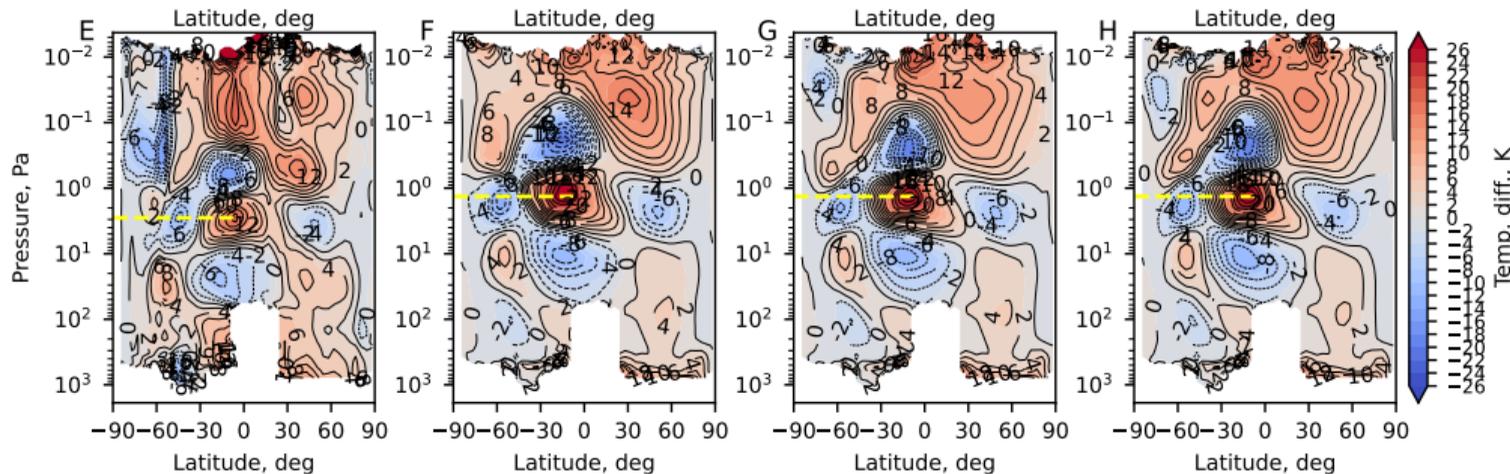
↑
dusty

Is 60-90, clear-sky

Temp.[K]

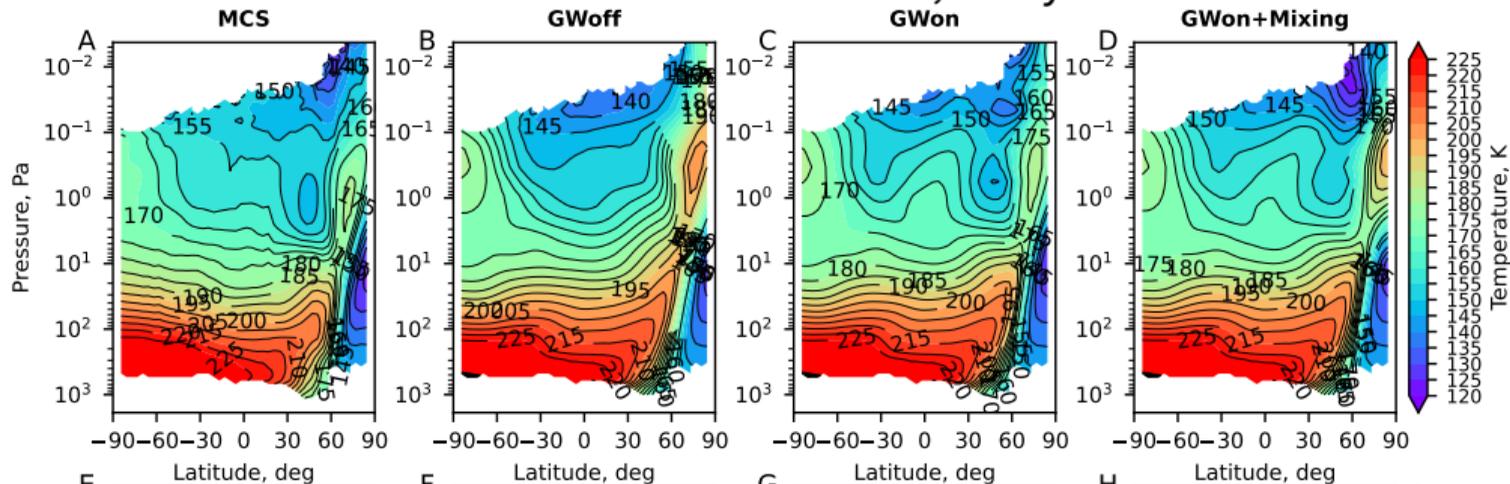


Tide [K]

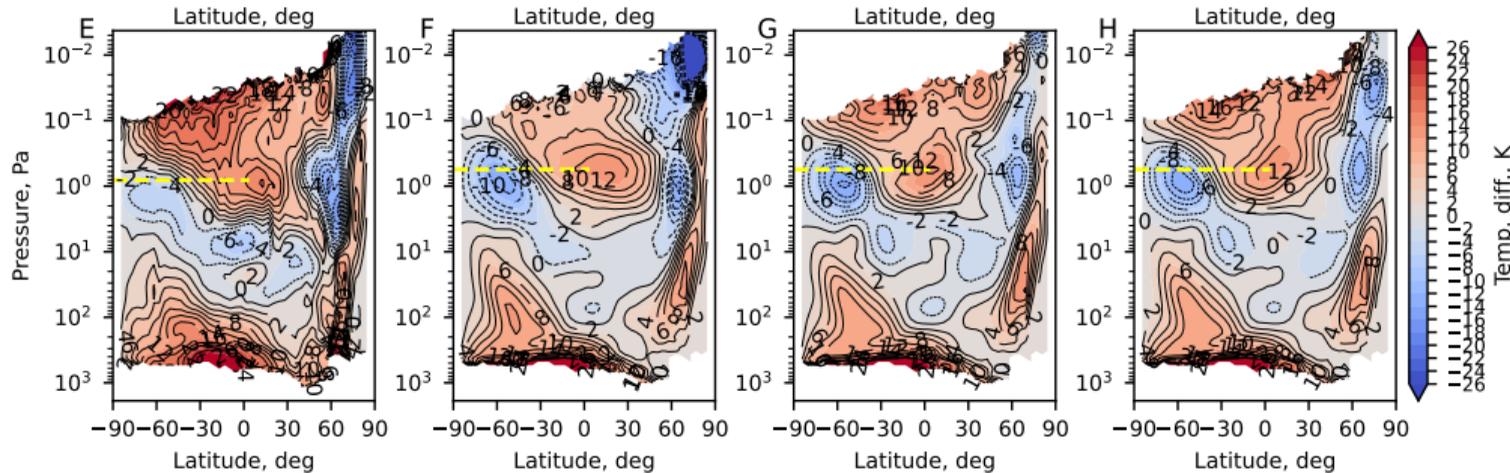


LS 240-270, dusty

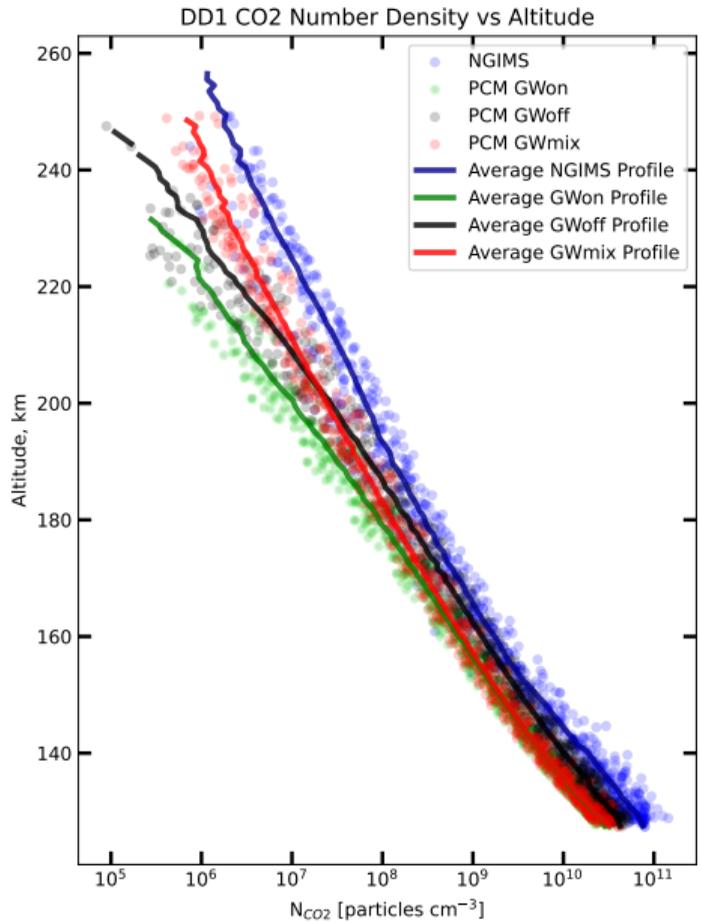
Temp. [K]



Tide [K]



MAVEN/NGIMS
GWoff
GWon
GWon+mixing



Thank you for your time and interest in this topic. Please feel free to ask any questions.

Je vous remercie pour votre temps et votre intérêt. N'hésitez pas à poser vos questions.

承蒙聆听，不胜荣幸，有疑但问。

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