

1 **Differential treatment of the PBL between inner and**
2 **outer regions of cold pools; I: over ocean**

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6 **Key Points:**

- 7 • cold pool strength is strongly reduced when the splitting of the PBL is taken into
8 account.
9 • cold pool strength is closer to observations when the splitting of the PBL is taken
10 into account.
11 • simulated precipitation variability is improved.

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12 **Abstract**
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14 **1 Introduction**

15 **2 Surface atmosphere coupling**

16 *2.0.0.1 Notations* : We consider fields function of position (x, y) and time t , where
 17 (x, y) belongs to a very large domain (large when compared to cold pool sizes and to grid
 18 cell size). Most often the t dependance will be omitted. For each field, say ϕ , ϕ^* desig-
 19 nates the average values over the domain, ϕ^w the average value over the (w) region, and
 20 ϕ^x the average value over the (x) region. Fluxes are positive downward.

21 At the surface, the boundary layer model is coupled with the subsurface model. The
 22 subsurface model may represent surface water of an ocean, soil at the surface of conti-
 23 nent, land ice, or sea ice. In all cases we assume that the subsurface model is constrained
 24 by mixed boundary conditions, that is by an affine relationship between the surface hu-
 25 midity q_s^* and the surface moisture flux $\hat{\phi}^*$:

$$\hat{\phi}^* = \hat{\mu} - \hat{\lambda}q_s^* \quad (1)$$

26 The coupling between the two models is implemented in the following way: the bound-
 27 ary layer model computes the coefficients $\hat{\lambda}$ and $\hat{\mu}$; from these boundary conditions the
 28 subsurface model determines the values of the variables $\hat{\phi}^*$ and q_s^* ; from these surface
 29 values the boundary layer scheme computes the humidity values in the whole troposphere.

30 However it is not directly $\hat{\lambda}$ and $\hat{\mu}$ that are used as boundary conditions for the sub-
 31 surface model but another set of equivalent parameters. An intermediate variable is in-
 32 troduced, namely the apparent atmosphere humidity q^a , related to the surface flux by:

33 1. a "surface exchange" like equation:

$$\hat{\phi}^* = \hat{K}^a(q^a - q_s^*) \quad (2)$$

34 2. the sensivity coefficients of the boundary layer model to the surface flux:

$$q^a = \hat{A}^a + \hat{B}^a\hat{\phi}^*\Delta t \quad (3)$$

35 Then the boundary conditions are given by the three coefficients \hat{K}^a , \hat{A}^a and \hat{B}^a . The
 36 equivalence with the boundary condition (1) may be shown by eliminating q^a between
 37 the two equations (2) and (3), which yields a relation of the form (1) with:

$$\begin{aligned} \hat{\mu} &= \frac{\hat{K}^a\hat{A}^a}{1 - \hat{K}^a\hat{B}^a\Delta t} \\ \hat{\lambda} &= \frac{\hat{K}^a}{1 - \hat{K}^a\hat{B}^a\Delta t} \end{aligned} \quad (4)$$

38 For a given pair $(\hat{\mu}, \hat{\lambda})$ there exists an infinity of triplets $(\hat{K}^a, \hat{A}^a, \hat{B}^a)$ such that
 39 the relations (4) hold. \hat{K}^a may be chosen arbitrarily different from zero and \hat{A}^a and \hat{B}^a
 40 are given by:

$$\begin{aligned} \hat{A}^a &= \frac{\hat{\mu}}{\hat{\lambda}} \\ \hat{B}^a &= \frac{1}{\Delta t} \left[\frac{1}{\hat{K}^a} - \frac{1}{\hat{\lambda}} \right] \end{aligned} \quad (5)$$

In the LMDZ GCM \hat{K}^a is chosen equal to the exchange coefficient between the first layer of the atmospheric model and the surface in which case: (i) the surface flux reads $\hat{\phi}^* = \hat{K}(q_1^* - q_s^*)$ so that q^a is identical to q_1^* , and (ii) the coefficients \hat{A}^a and \hat{B}^a are directly given by the resolution of the vertical diffusion equation in the atmosphere.

It is worth emphasizing that the equality $\hat{K}^a = \hat{K}$ (and its consequence $q^a = q_1^*$) is chosen in order to make it possible to separate the part of the subsurface boundary conditions due to surface processes (\hat{K}^a) and the part due to boundary layer processes (\hat{A}^a and \hat{B}^a). When dealing with a split boundary layer, even though the same consideration of physical significance is accounted for, q^a will be different from q_1^* .

3 Splitting

For the sake of brevity, we present the full computations only for humidity. We shall outline the differences concerning enthalpy afterward.

3.1 Basic equations

The main assumptions of the model are:

1. All fields are horizontally homogeneous within (w) and within (x).
2. Model equations are linear during each time step: exchange coefficients are supposed constant (they are computed with the field values at the end of the previous time step).
3. Surface fluxes are given by:

$$\begin{cases} \hat{\phi}^w &= \hat{K}^w(q_1^w - q_s^w) \\ \hat{\phi}^x &= \hat{K}^x(q_1^x - q_s^x) \end{cases} \quad (6)$$

4. Surface moistures q_s^w and q_s^x differ by a prescribed amount $\hat{\Delta}_s$:

$$\delta q_s = \hat{\Delta}_s \quad (7)$$

5. The boundary layer scheme provides, within each region (w) and (x), a linear relationship between humidity q_1 at the first level and the surface moisture flux $\hat{\phi}$:

$$\begin{cases} q_1^w &= (\hat{A}^w + \hat{B}^w \hat{\phi}^w \Delta t) \\ q_1^x &= (\hat{A}^x + \hat{B}^x \hat{\phi}^x \Delta t) \end{cases} \quad (8)$$

In addition to these five equations, (6), (7) and (8), there are six equations relating for each of the field $\hat{\phi}$, q_s and q_1 the average value, the values in the regions (w) and (x), and the difference between the values in the regions (w) and (x). For instance for the moisture flux:

$$\begin{cases} \hat{\phi}^* &= \sigma_w \hat{\phi}^w + (1 - \sigma_w) \hat{\phi}^x \\ \delta \hat{\phi} &= \hat{\phi}^w - \hat{\phi}^x \end{cases} \quad (9)$$

and similarly for the surface moisture difference and for the humidity at the first level.

To sum up, there are eleven linear equations for the twelve variables $\hat{\phi}^*$, $\hat{\phi}^w$, $\hat{\phi}^x$, $\delta \hat{\phi}$, q_1^* , q_1^w , q_1^x , δq_1 , q_s^* , q_s^w , q_s^x , δq_s . from which it is possible (by elimination of all variables but $\hat{\phi}^*$ and q_s^*) to extract an affine relationship between the average flux $\hat{\phi}^*$ and the average surface humidity q_s^* :

$$\hat{\phi}^* = \hat{\mu} - \hat{\lambda} q_s^* \quad (10)$$

where $\hat{\mu}$ and $\hat{\lambda}$ are coefficients which may be expressed in terms of the coefficients of the eleven linear relations. It is this relation which constitutes the mixed boundary condition for the surface model.

75 The purpose of the following is to determine the coefficients $\hat{\mu}$ and $\hat{\lambda}$. Moreover,
 76 in order to facilitate the coupling with surface models, we shall rewrite the mixed bound-
 77 ary conditions in the form of equations similar to equations (6) and (8): we shall first
 78 determine \hat{A}^a and \hat{B}^a assuming that \hat{K}^a is known. Then we shall choose \hat{K}^a in such a
 79 way that it is only function of surface conditions (and not of the boundary layer above).

80 3.2 Solving

81 3.2.1 Determining $\hat{\lambda}$ and $\hat{\mu}$

82 We introduce the coefficients \hat{K}' :

$$\left\{ \begin{array}{l} \hat{K}'^x = \frac{\hat{K}^x}{1 - \hat{K}^x \hat{B}^x \Delta t} \\ \hat{K}'^w = \frac{\hat{K}^w}{1 - \hat{K}^w \hat{B}^w \Delta t} \end{array} \right. \quad (11)$$

83 These coefficients are effective exchange coefficients taking into account the boundary
 84 layer feedbacks:

$$\begin{aligned} \hat{\phi}^w &= \hat{K}'^w (\hat{A}^w - q_s^w) \\ \hat{\phi}^x &= \hat{K}'^x (\hat{A}^x - q_s^x) \end{aligned} \quad (12)$$

85 Now applying the second product identity (A.4) to the fields \hat{K}' , $\hat{A} - q_s$ and their prod-
 86 uct $\hat{\phi}$ yields:

$$\hat{\phi}^* = \hat{K}'^* (\hat{A}^* - q_s^*) + \sigma_w \sigma_x \delta \hat{K}' (\delta \hat{A} - \delta q_s) \quad (13)$$

87 or, since δq_s is prescribed (Eq. 7):

$$\hat{\phi}^* = \hat{K}'^* \hat{A}^* + \sigma_w \sigma_x \delta \hat{K}' (\delta \hat{A} - \hat{\Delta}_s) - \hat{K}'^* q_s^* \quad (14)$$

88 We recognize in this equation the mixed boundary condition we are seeking. Identifica-
 89 tion with equation (1) yields:

$$\begin{aligned} \hat{\mu} &= \hat{K}'^* \hat{A}^* + \sigma_w \sigma_x \delta \hat{K}' (\delta \hat{A} - \hat{\Delta}_s) \\ \hat{\lambda} &= \hat{K}'^* \end{aligned} \quad (15)$$

90 Now we want to write the boundary conditions in terms of the triplet $(\hat{K}^a, \hat{A}^a, \hat{B}^a)$.
 91 As explained in section (2), since \hat{K}^a is then a free parameter, we first determine \hat{A}^a and
 92 \hat{B}^a for a given \hat{K}^a .

93 3.2.2 Solving for \hat{A}^a and \hat{B}^a for a given \hat{K}^a

94 The coefficients \hat{A}^a and \hat{B}^a are given by equations (5) where $\hat{\lambda}$ and $\hat{\mu}$ are given by
 95 equations (15):

$$\begin{aligned} \hat{A}^a &= \hat{A}^* + \sigma_w \sigma_x \frac{\delta \hat{K}'}{\hat{K}'^*} (\delta \hat{A} - \hat{\Delta}_s) \\ \hat{B}^a &= \frac{1}{\Delta t} \left[\frac{1}{\hat{K}^a} - \frac{1}{\hat{K}'^*} \right] \end{aligned} \quad (16)$$

The problem of coupling a split boundary layer with a uniform subsurface model is now solved. Whatever the non-zero coefficient \hat{K}^a the apparent atmosphere humidity reads

$$q^a = \hat{A}^a + \hat{B}^a \hat{\phi}^* \Delta t$$

A possible choice for \hat{K}^a would be $\hat{K}^a = \hat{K}'^*$. Then $\hat{B}^a = 0$ and $q^a = \hat{A}^a$ so that the mixed boundary condition for the subsurface model reads:

$$\hat{\phi}^* = \hat{K}^a (\hat{A}^a - q_s^*)$$

96 which means that the subsurface model behaves as if coupled to a fixed moisture atmo-
 97 sphere. This is a particularly simple formulation. However it has a drawback: the co-
 98 efficient \hat{K}^a is dependent on both the surface exchange processes and on the boundary
 99 layer processes. Moreover, obviously q^a (which does not depend on $\hat{\phi}^*$) will be in gen-
 100 eral different from q_1^* (which varies with $\hat{\phi}^*$), even if there is no cold pool. We shall now
 101 seek a more satisfactory formulation.

102 3.2.3 Determining \hat{K}^a

Similarly to the non-split case we might choose $\hat{K}^a = \hat{K}^*$. Although this will
 be our final choice, it needs some more justification, since the weighting between \hat{K}^w and
 \hat{K}^x could be very different from the one given by \hat{K}^* . In order to guide our choice we
 shall look at the variable q_1^* and compare it to q^a . The moisture field q_1 is related to the
 surface fluxes by equations (6):

$$\begin{cases} \hat{\phi}^w &= \hat{K}^w(q_1^w - q_s^w) \\ \hat{\phi}^x &= \hat{K}^x(q_1^x - q_s^x) \end{cases}$$

103 Using the same technique as in section (3.2.1) we apply the second product identity (A.4)
 104 to the fields \hat{K} , $q_1 - q_s$ and their product $\hat{\phi}$. It yields:

$$\hat{\phi}^* = \hat{K}^*(q_1^* - q_s^*) + \sigma_w \sigma_x \delta \hat{K} (\delta q_1 - \hat{\Delta}_s) \quad (17)$$

105 (It should be noted that this equation is not a mixed boundary condition for the sub-
 106 surface model since it involves four unknown variables: $\hat{\phi}^*$, q_1^* , q_s^* , and δq_1).

107 Elimination of $\hat{\phi}^*$ between equations (2) and (17) yields an expression for the dif-
 108 ference $q^a - q_1^*$:

$$q^a - q_1^* = \left(1 - \frac{\hat{K}^a}{\hat{K}^*}\right)(q^a - q_s^*) + \sigma_w \sigma_x \frac{\delta \hat{K}}{\hat{K}^*} (\delta q_1 - \hat{\Delta}_s) \quad (18)$$

This equation shows that when there are no cold pools ($\sigma_w = 0$ or $\sigma_x = 0$) then:

$$q^a - q_1^* = \left(1 - \frac{\hat{K}^a}{\hat{K}^*}\right)(q^a - q_s^*)$$

109 which implies that $q^a = q_1^*$ if and only if $\hat{K}^a = \hat{K}^*$. Consequently, in order to
 110 guarantee consistency with the no-cold pool case we choose $\hat{K}^a = \hat{K}^*$.

On the other hand, the equation for \hat{B}^a may be rewritten in the form of a sum in-
 stead of a difference. First \hat{B}^a reads also $\hat{B}^a = (1/\Delta t)(\hat{K}'^* - \hat{K}^*)/(\hat{K}^* \hat{K}'^*)$. Sec-
 ond it is easy to show that:

$$\hat{K}'^* - \hat{K}^* = \Delta t [\sigma_w \hat{K}^w \hat{K}'^w \hat{B}^w + \sigma_x \hat{K}^x \hat{K}'^x \hat{B}^x]$$

111 Hence \hat{B}^a reads also:

$$\hat{B}^a = \frac{\sigma_w \hat{K}^w \hat{K}'^w \hat{B}^w + \sigma_x \hat{K}^x \hat{K}'^x \hat{B}^x}{\hat{K}^* \hat{K}'^*} \quad (19)$$

112 Hence the final formulas:

$$\begin{cases} \hat{K}^a &= \hat{K}^* \\ \hat{A}^a &= \hat{A}^* + \sigma_w \sigma_x \frac{\delta \hat{K}'^*}{\hat{K}'^*} (\delta \hat{A} - \hat{\Delta}_s) \\ \hat{B}^a &= \frac{\sigma_w \hat{K}^w \hat{K}'^w \hat{B}^w + \sigma_x \hat{K}^x \hat{K}'^x \hat{B}^x}{\hat{K}^* \hat{K}'^*} \end{cases} \quad (20)$$

113 Equations (20) express the boundary conditions for the subsurface model coupled
 114 to an equivalent horizontally homogeneous atmosphere. It enables the subsurface model
 115 to determine the average surface flux $\hat{\phi}^*$. There remain to determine the repartition of
 116 this moisture flux among the two regions (w) and (x); otherwise stated, there remain to
 117 determine the difference $\delta\hat{\phi}$.

118 3.3 Back to the atmosphere

119 In order to compute the flux difference $\delta\hat{\phi}$ we apply the third product identity (A.6)
 120 to the fields \hat{K}' , $\hat{A} - q_s$ and their product $\hat{\phi}$:

$$\hat{K}'^* \delta\hat{\phi} - \delta\hat{K}' \hat{\phi}^* = \hat{K}'^x \hat{K}'^w (\delta\hat{A} - \delta q_s) \quad (21)$$

121 Hence the expression for $\delta\hat{\phi}$:

$$\delta\hat{\phi} = \frac{\delta\hat{K}'}{\hat{K}'^*} \hat{\phi}^* + (\delta\hat{A} - \Delta_s) \frac{\hat{K}'^x \hat{K}'^w}{\hat{K}'^*} \quad (22)$$

122 3.4 Formulas for static energy

123 The equations for the dry static energy $h = C_p T + gz$ are almost identical to
 124 those for humidity except for the surface variable: it is q_s in the case of humidity while
 125 it is $T_s = h_s/C_p$ in the case of dry static energy. Hence the equations equivalent to
 126 (15) read:

$$\mu = K'^* A^* + \sigma_w \sigma_x \delta K' (\delta A - C_p \Delta_s) \quad (23)$$

$$\lambda = K'^*$$

127 , the equations equivalent to (2) and (3) read:

$$\begin{cases} \phi^* = K^a (h^a - C_p T_s^*) \\ h^a = A^a + B^a \phi^* \Delta t \end{cases} \quad (24)$$

128 , and the equations equivalent to (20) read:

$$\begin{cases} K^a = K^* \\ A^a = A^* + \sigma_w \sigma_x \frac{\delta K'}{K'^*} (\delta A - C_p \Delta_s) \\ B^a = \frac{\sigma_w K^w K'^w B^w + \sigma_x K^x K'^x B^x}{K^* K'^*} \end{cases} \quad (25)$$

129 Finally the "return to atmosphere" formula reads:

$$\delta\hat{\phi} = \frac{\delta K'}{K'^*} \hat{\phi}^* + (\delta A - C_p \Delta_s) \frac{K'^x K'^w}{K'^*} \quad (26)$$

130 4 1D simulations

131 5 3D simulations

132 6 Conclusion

133 A: Elementary product identities

134 Let a , b and p be three fields such that $p^w = a^w b^w$ and $p^x = a^x b^x$. This ap-
 135 pendix proves three identities relating the field a , b and their product p .

136 Since $a^w = a^* + \sigma_x \delta a$ and $a^x = a^* - \sigma_w \delta a$ (and similarly for b), the products
 137 p^w and p^x read:

$$\begin{aligned} p^w &= (a^* + \sigma_x \delta a)(b^* + \sigma_x \delta b) \\ p^x &= (a^* - \sigma_w \delta a)(b^* - \sigma_w \delta b) \end{aligned} \quad (\text{A.1})$$

138 Hence:

$$\begin{aligned} p^w &= a^* b^* + \sigma_x^2 \delta a \delta b + \sigma_x (a^* \delta b + b^* \delta a) \\ p^x &= a^* b^* + \sigma_w^2 \delta a \delta b - \sigma_w (a^* \delta b + b^* \delta a) \end{aligned} \quad (\text{A.2})$$

139 From which it is possible to derive expressions for δp ($= p^w - p^x$) and p^* ($= \sigma_w p^w +$
 140 $\sigma_x p^x$) yielding the first product identity:

$$\delta p = a^* \delta b + b^* \delta a + (\sigma_x - \sigma_w) \delta a \delta b \quad (\text{A.3})$$

141 and the second product identity:

$$p^* = a^* b^* + \sigma_w \sigma_x \delta a \delta b \quad (\text{A.4})$$

142 Another usefull identity is obtained by forming $a^* \delta p - p^* \delta a$:

$$\begin{aligned} a^* \delta p - p^* \delta a &= \delta b [(a^*)^2 + (\sigma_x - \sigma_w) a^* \delta a - \sigma_x \sigma_w (\delta a)^2] \\ &= \delta b (a^* + \sigma_x \delta a)(a^* - \sigma_w \delta a) \end{aligned} \quad (\text{A.5})$$

143 Finally the third product identity reads:

$$a^* \delta p - p^* \delta a = a^x a^w \delta b \quad (\text{A.6})$$